

Power exhaust in ITER II: Divertor

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with grateful acknowledgement for the contributions of

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Content

- Introduction to the ITER W divertor
 - Basic physics/design features and expected lifetime
- Stationary power loading – the design simulation database
 - Overall characteristics
 - Focus on factors influencing the peak power loading
 - What really is the tolerable steady state power flux density?
 - What are the tolerable ELM loads?
- I will speak only about axisymmetric divertor heat loads → see lecture by O. Schmitz for the case of 3D fields for ELM control

But first, a brief return to lecture 1

- Apply 0D power balance to estimate peak q_{\parallel} at the divertor targets, e.g. for outer target:

$$q_{\parallel out} = P_{div,out} / (2\pi R_{out} \lambda_q (B_{\theta} / B_{\phi})_{omp}) \quad P_{div,out} = P_{SOL} (1 - f_{RAD}) A_{sym} / (1 + A_{sym})$$

f_{RAD} = radiated fraction of power conducted to the divertor

R_{out} = major radius of outer strike point

$P_{div,out}$, $P_{div,in}$ = powers into outer/inner divertor, $A_{sym} = P_{div,out} / P_{div,in}$

- Example: $\lambda_q = 5$ mm, $P_{SOL} \sim 100$ MW, $A_{sym} = 2$
“detached divertor”, $f_{RAD} \sim 60\% \rightarrow q_{\parallel out} \sim 300$ MWm⁻²
“low recycling divertor”, $f_{RAD} \sim 20\%$, $q_{\parallel out} \sim 900$ MWm⁻²
- Take $\alpha \sim 4^\circ$ (component shaping \rightarrow see later) $\rightarrow q_{\perp,out} = q_{\parallel,out} \sin\alpha$
 $q_{\perp,out} \sim 21$ MWm⁻² (“detached”), $q_{\perp,out} \sim 63$ MWm⁻² (“attached”)!!
- Cannot be handled by technology on ITER

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NB: no losses associated with diffusion into PFR included \rightarrow see lecture by H. Zohm

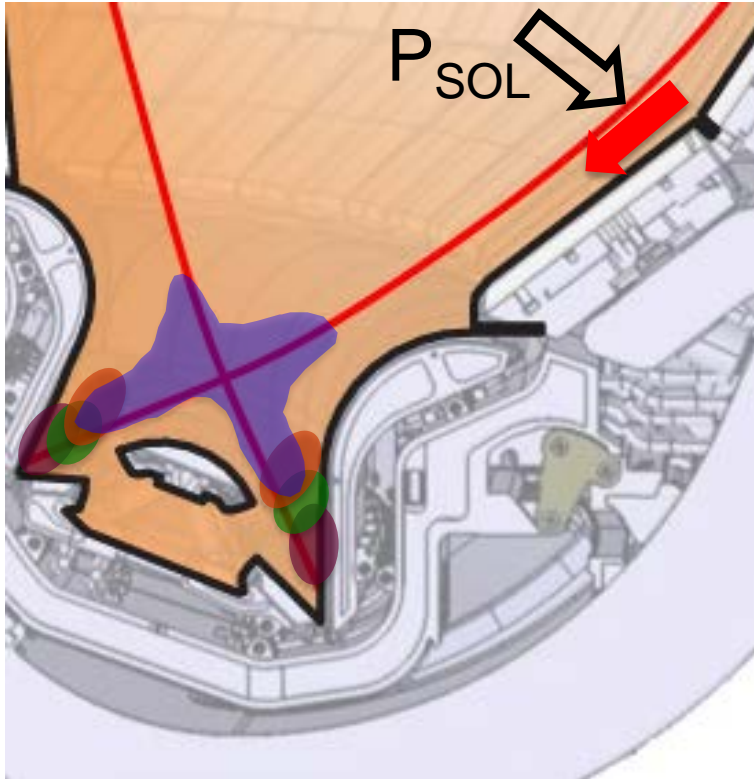
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Now imagine that λ_q drops by factor 5 $\rightarrow q_{\perp,out}$ increases by same factor ...

- Cannot be handled by technology on ITER

Problem is that simple specification is too simple

- ITER Divertor is highly dissipative



Heat conduction zone

Impurity radiation zone

**$H^0/D^0/T^0$ ionization zone
($T_e > 5$ eV)**

Neutral friction zone

**Recombination zone
($T_e < 1$ eV)**

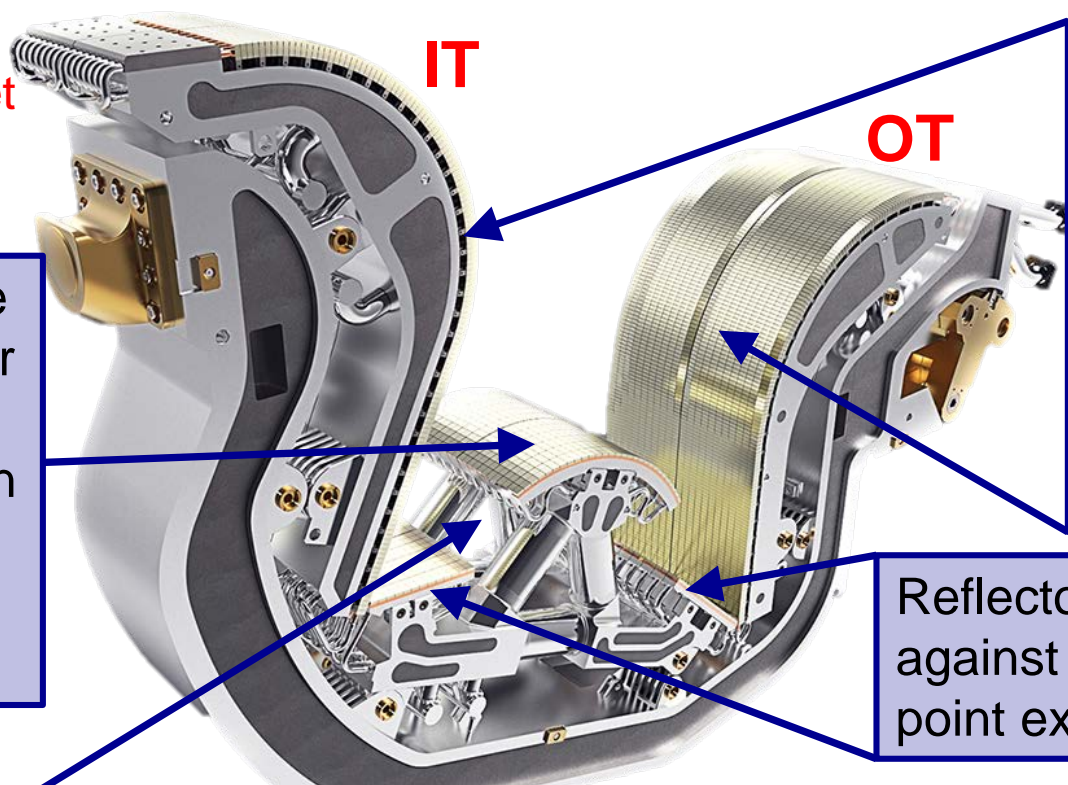
The ITER tungsten divertor



- The most sophisticated tokamak divertor ever built
 - 54 individual cassettes, fully water cooled, designed to handle up to ~100 MW in steady state
 - Now entering the procurement phase → design essentially complete

W divertor: key physics characteristics

IT = INNER target
OT = OUTER target



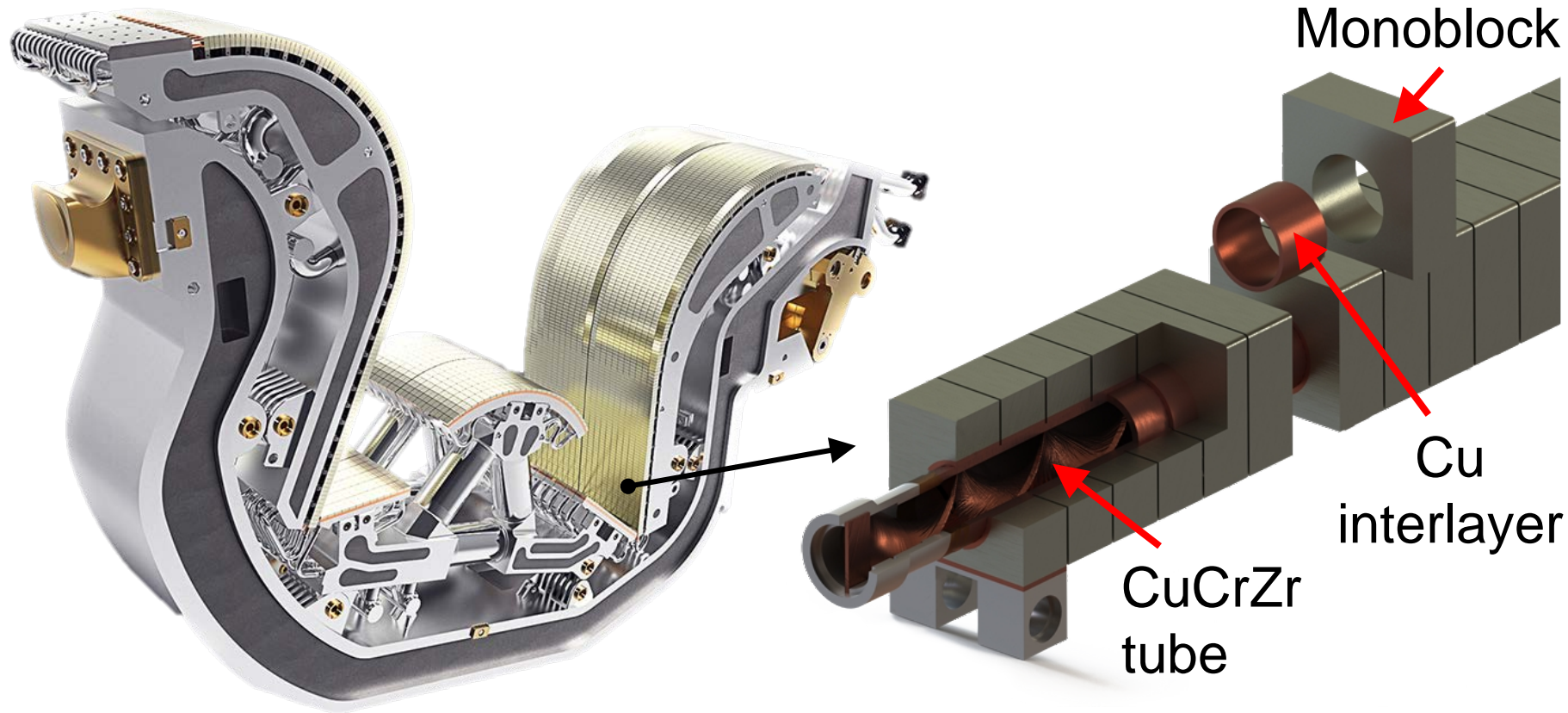
Dome to improve pumping \rightarrow lower pumping speed required for given upstream He conc or fuel throughput

Deep vertical targets and baffle regions promoting detachment and reducing neutral escape to the core

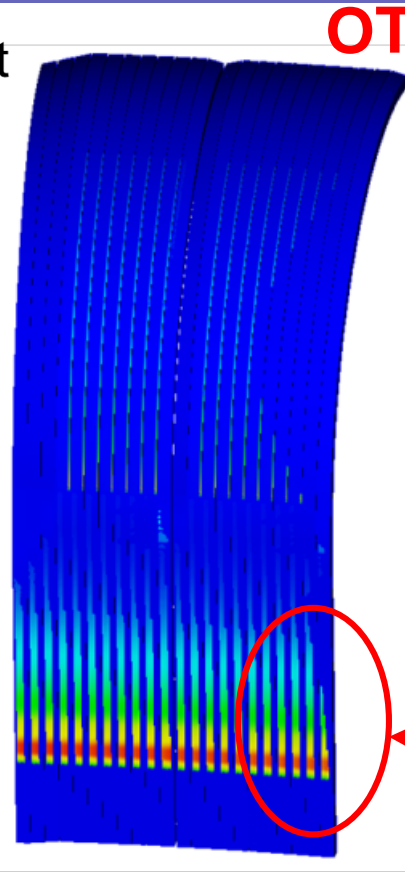
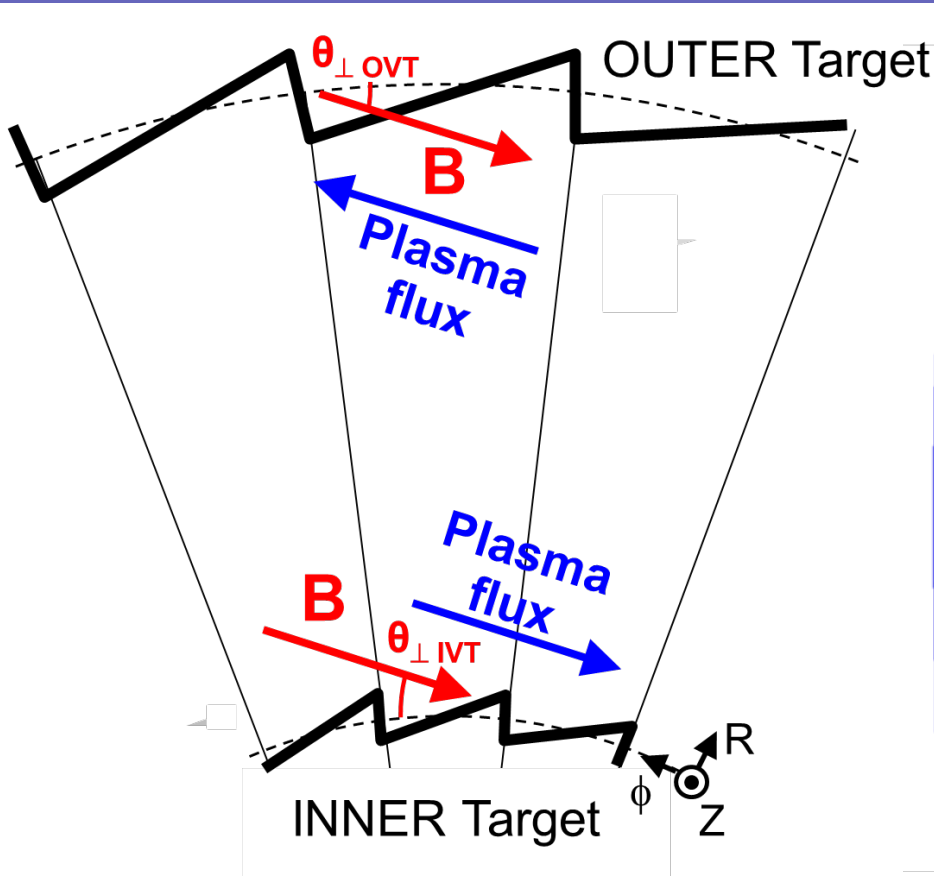
Reflector plates protect against downward strike point excursions

Transparency between targets for neutral recirculation – lower power asymmetries

Vertical target plasma-facing units



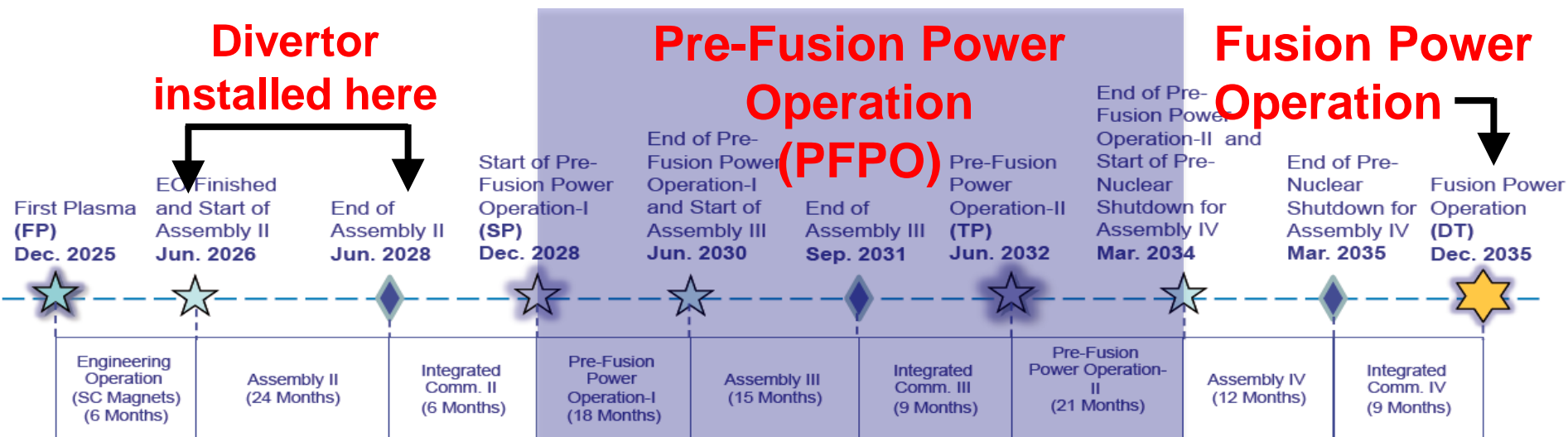
Global shaping: target tilting



- Tilt of $\sim 0.5^\circ$ at both inner and outer targets to protect gross leading edges between adjacent cassettes

First plasma-facing unit fully magnetically shadowed

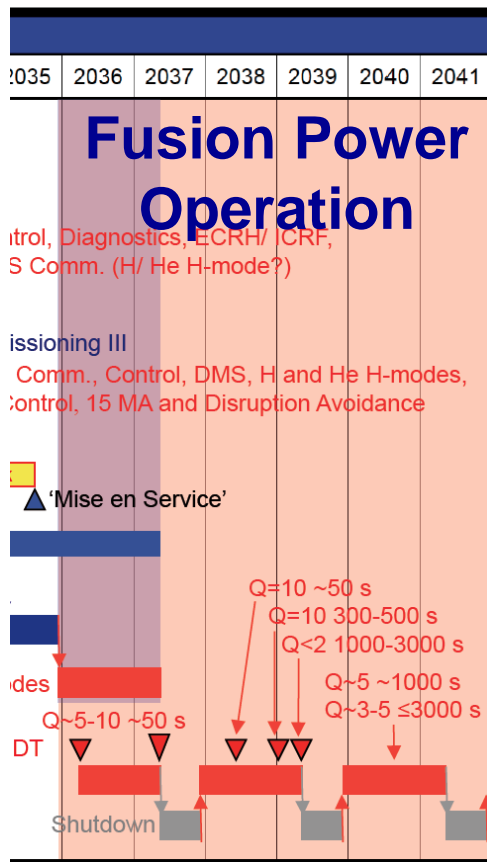
Revised ITER schedule and divertor lifetime



- H and He operation in PFPO phases 1 and 2, L- and H-mode → detailed operational breakdown in new ITER Research Plan*:
 - Total days in PFPO-1: 470 → ~5700 pulses → $\sim 3 \times 10^5$ s up to $P_{\text{heat}} = 30$ MW
 - Total days in PFPO-2: 545 → ~5600 pulses → $\sim 6 \times 10^5$ s up to $P_{\text{heat}} = 73$ MW

* <https://www.iter.org/technical-reports>

Revised ITER schedule and divertor lifetime



- First DT campaigns roughly split into 3 phases (FPO-1,2,3):
 - Expect power into SOL to reach design value ~100 MW
 - Expect ~900 days operation over ~5 years
 - ~12,000 pulses
 - **~8x10⁶ s plasma time** (~2200 hours or ~90 days)

First ITER W divertor required to survive until end of FPO-3

Divertor operation: design by simulation

- Physics operating mode for the ITER divertor is to a large extent based on plasma boundary simulations conducted over ~15 years with the SOLPS-4.3 code (B2-Eirene) → mostly C targets!

- Well documented

A. S. Kukushkin et al. J. Nucl. Mat. **290-293** (2001) 887
A. S. Kukushkin et al. Nucl. Fusion **42** (2002) 187
A. S. Kukushkin and H. D. Pacher, PPCF **44** (2002) 931
H. D. Pacher et al. J. Nucl. Mat. **313-316** (2003) 657
A. S. Kukushkin et al. Nucl. Fusion **43** (2003) 716
A. S. Kukushkin et al. Fus. Eng. Design **65** (2003) 355
A. S. Kukushkin et al. Nucl. Fusion **45** (2005) 608
A. S. Kukushkin et al. J. Nucl. Mat. **337-339** (2005) 17
A. S. Kukushkin et al. Nucl. Fusion **47** (2007) 698
A. S. Kukushkin et al. J. Nucl. Mat. **363-365** (2007) 308
G. W. Pacher et al. Nucl. Fusion **48** (2008) 105003
H. D. Pacher et al. J. Nucl. Mat. **390-391** (2009) 259
A. S. Kukushkin et al. Nucl. Fusion **49** (2009) 075008
A. S. Kukushkin et al., J. Nucl. Mat. **415** (2011) 2011
H. D. Pacher et al. J. Nucl. Mat. **415** (2011) S492
G. W. Pacher et al. Nucl. Fusion **51** (2011) 083004
A. S. Kukushkin et al. Fus. Eng. Design **86** (2011) 2865

A. S. Kukushkin et al. Nucl. Fusion **53** (2013) 123024

H. D. Pacher et al. J. Nucl. Mat. **463** (2015) 591

A. S. Kukushkin et al., Nucl. Fusion **56** (2016) 126012

First real operating domain study for metal walls → will be a focus of this talk

Since 2015, moved to new code version (incl. drifts)

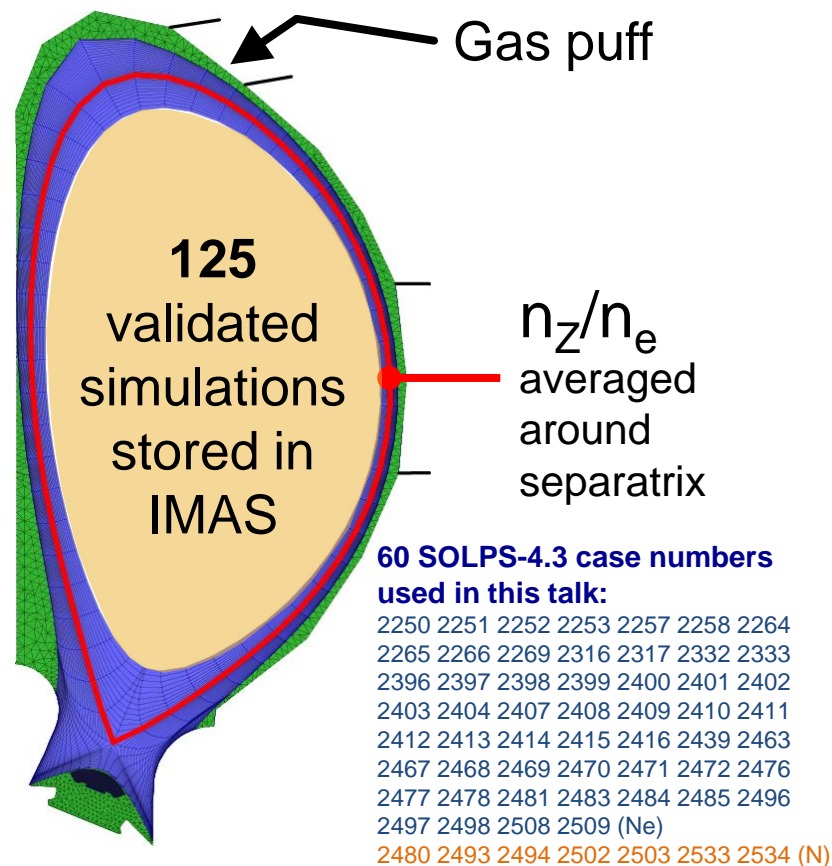
SOLPS-ITER

S. Wiesen et al., J. Nucl. Mat. **463** (2015) 480

X. Bonnin et al., Plasma & Fusion Research, **11** (2016) 1403102

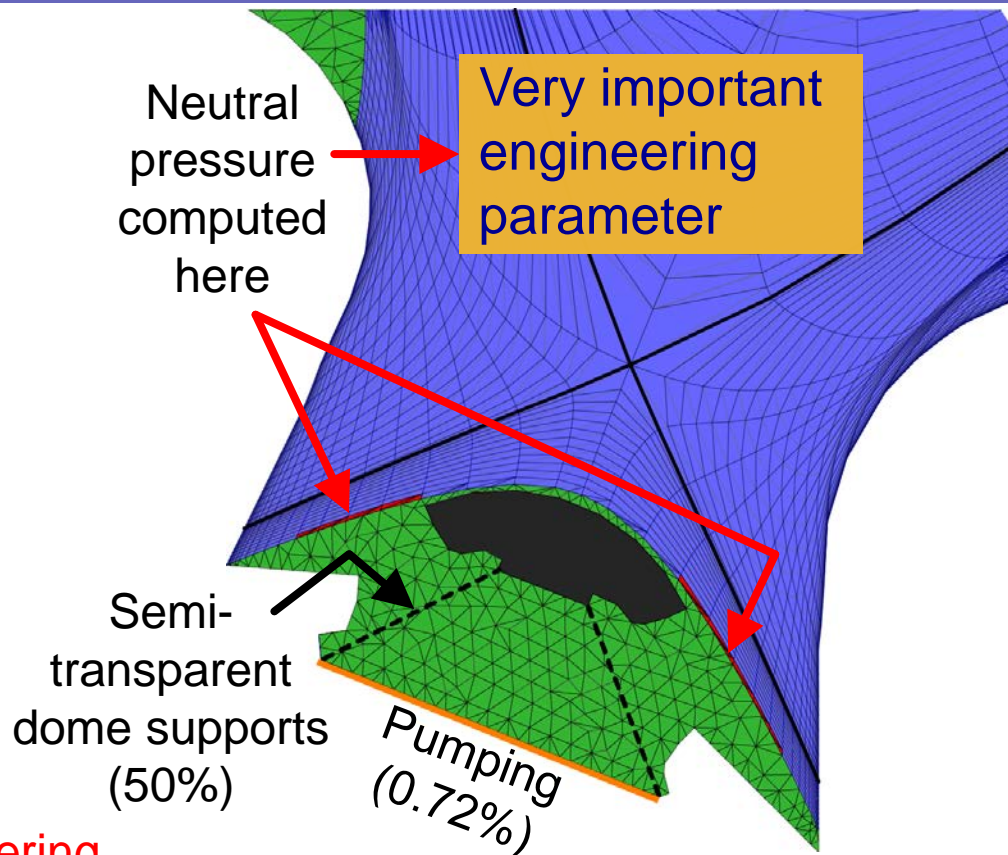
Main simulation database parameters

- Steady state – no ELMs
- No fluid drifts, “L-mode” edge
 - Neutral-neutral collisions included
- Fixed equilibrium
 - $q_{95} = 3$, $B_T/I_p = 1.8/5, 2.65/7.5, 5.3/15$
- Fixed cross-field transport
 - $D_{\perp} = 0.3 \text{ m}^2\text{s}^{-1}$, $\chi_{\perp} = 1.0 \text{ m}^2\text{s}^{-1}$
- Scans in fueling, seed impurity, power into numerical grid (P_{IN})
 - H, He, D, N_2 , Ne, but only $P_{IN} = 100 \text{ MW}$ in this talk
- All-metal walls
 - Assume Be everywhere, but no sputtering



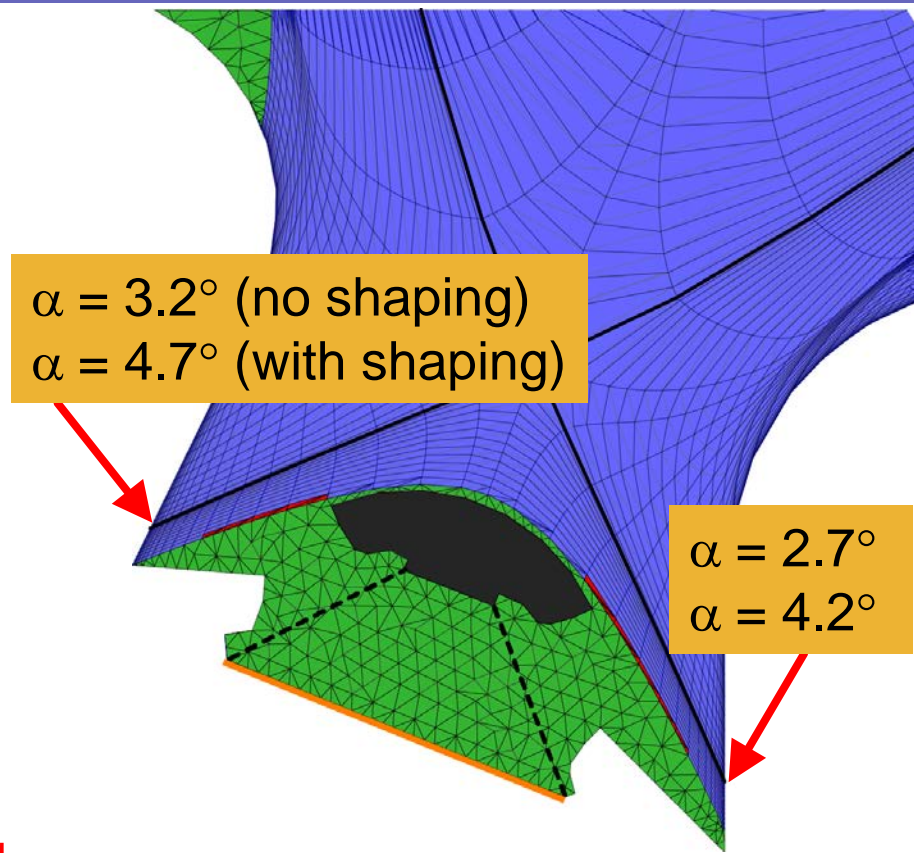
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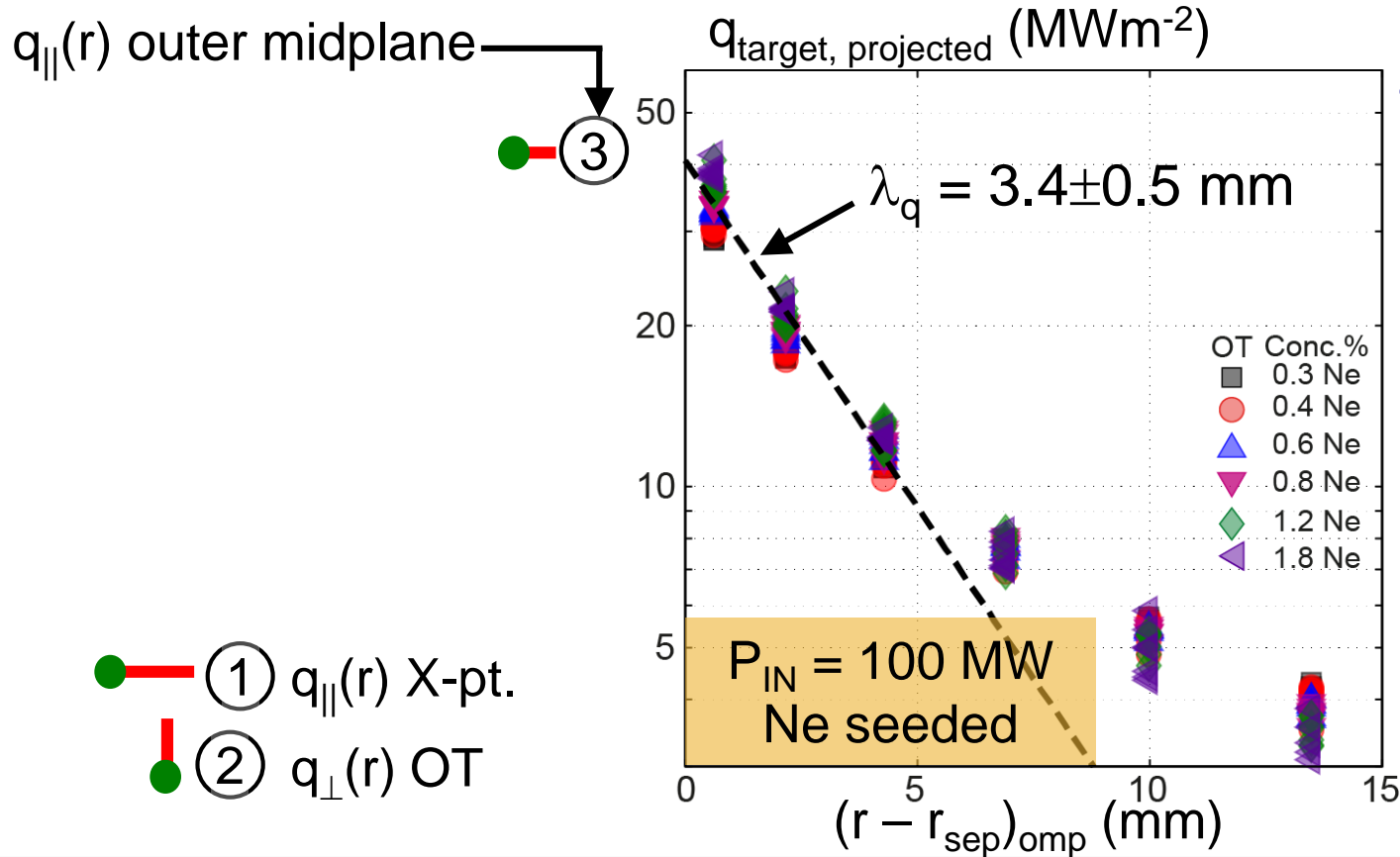
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Burning plasma operating window

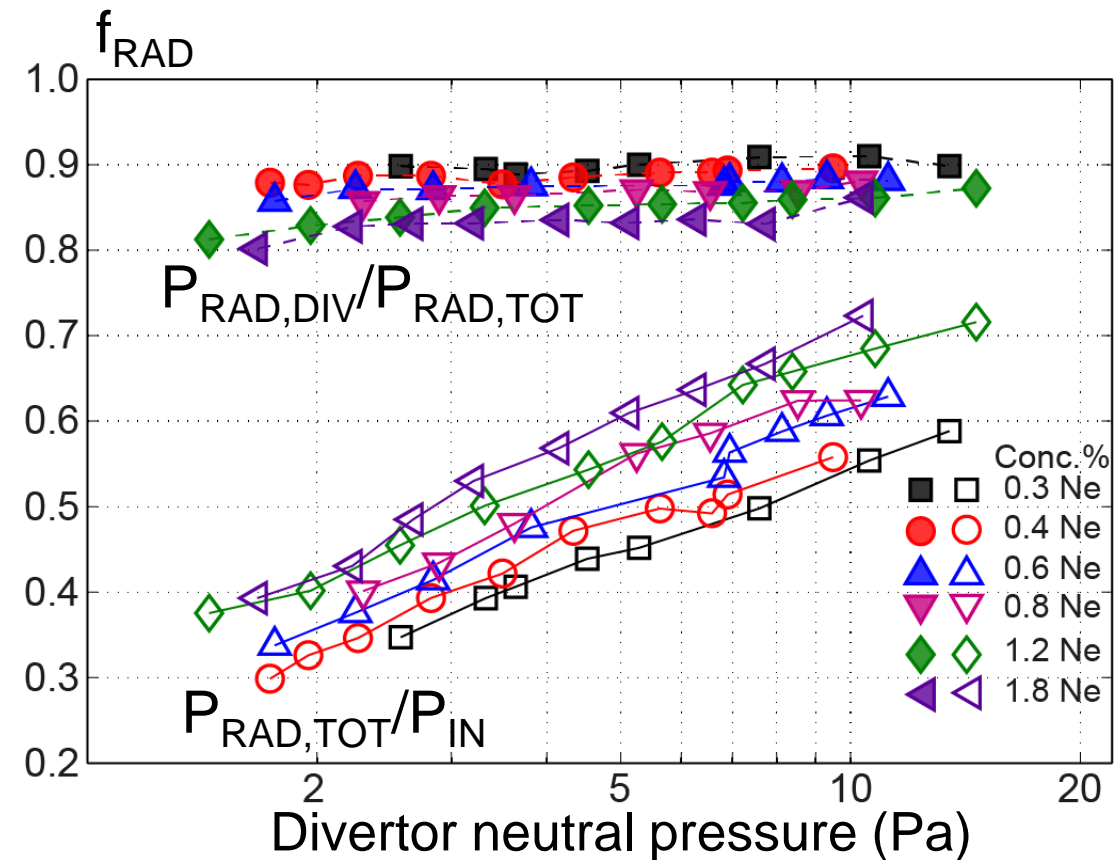
- Focus on “burning plasma” conditions → the most challenging for the ITER divertor
 - $Q_{DT} = 10$, $P_{IN} \sim 100$ MW
 - Ne and N₂ seeding (emphasis on Ne where database currently largest)
 - No discussion of “integrated modelling” here
- An important fact to bear in mind: ITER will operate always quite close to the H-mode power transition threshold
 - Cannot afford (too) much edge/core radiation

SOL heat flux width



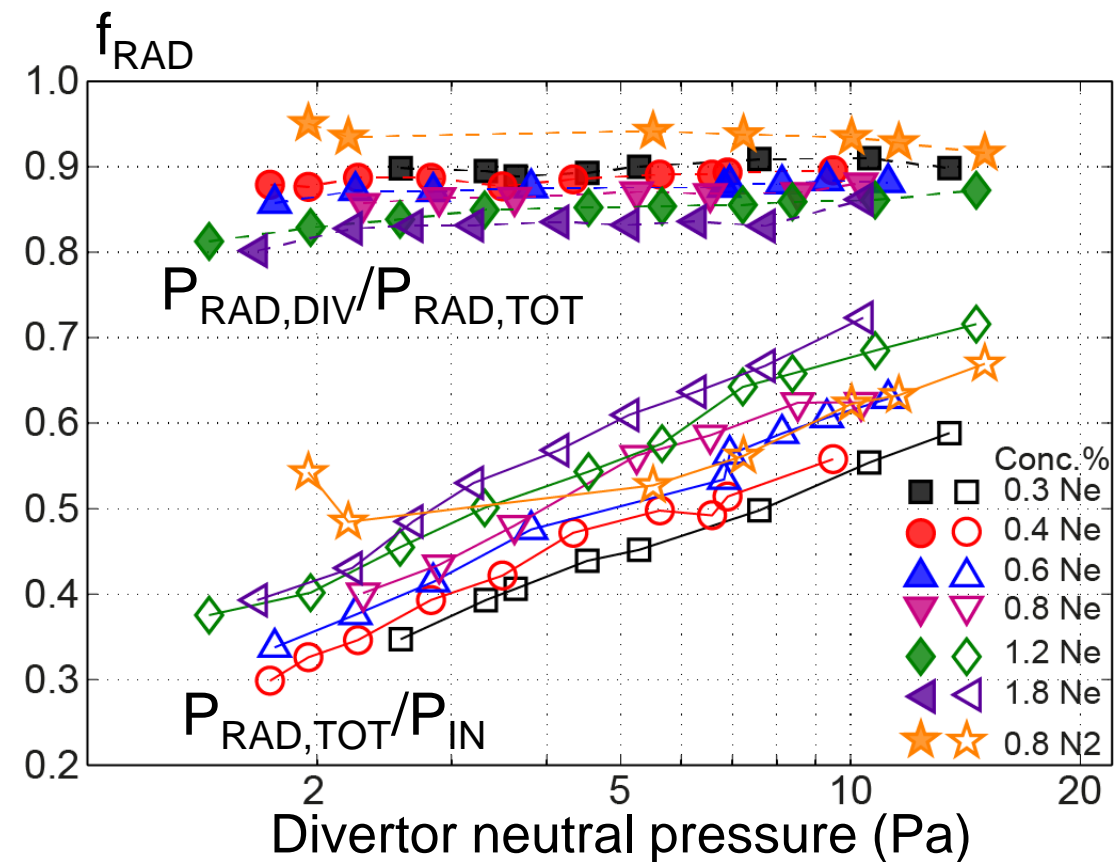
- Divertor conditions across database do not strongly influence upstream λ_q

Radiated fractions



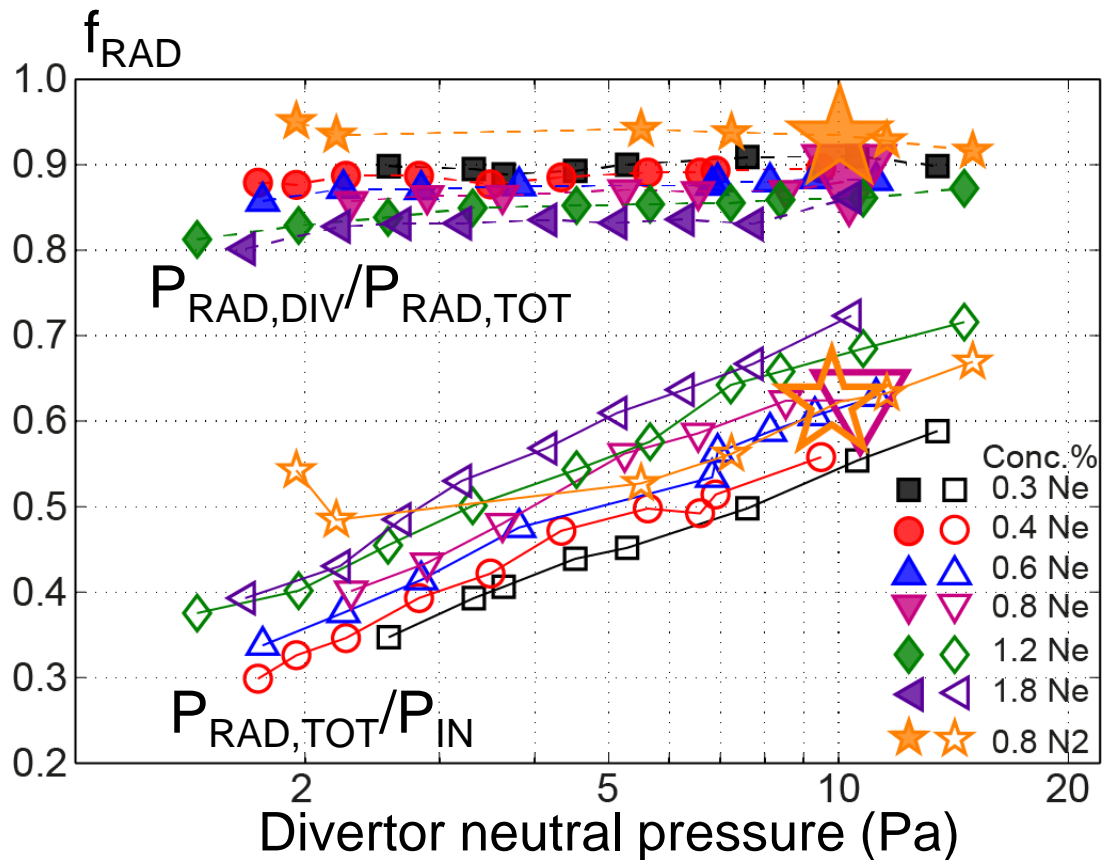
- Radiation largely confined to the divertor region
 - $f_{RAD,DIV} \sim 0.8-0.9$ across operating window for Ne
 - $f_{RAD,TOT} \sim 0.3 - 0.7$

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 - N more efficiently compressed than Ne
 - Lower core radiation with N

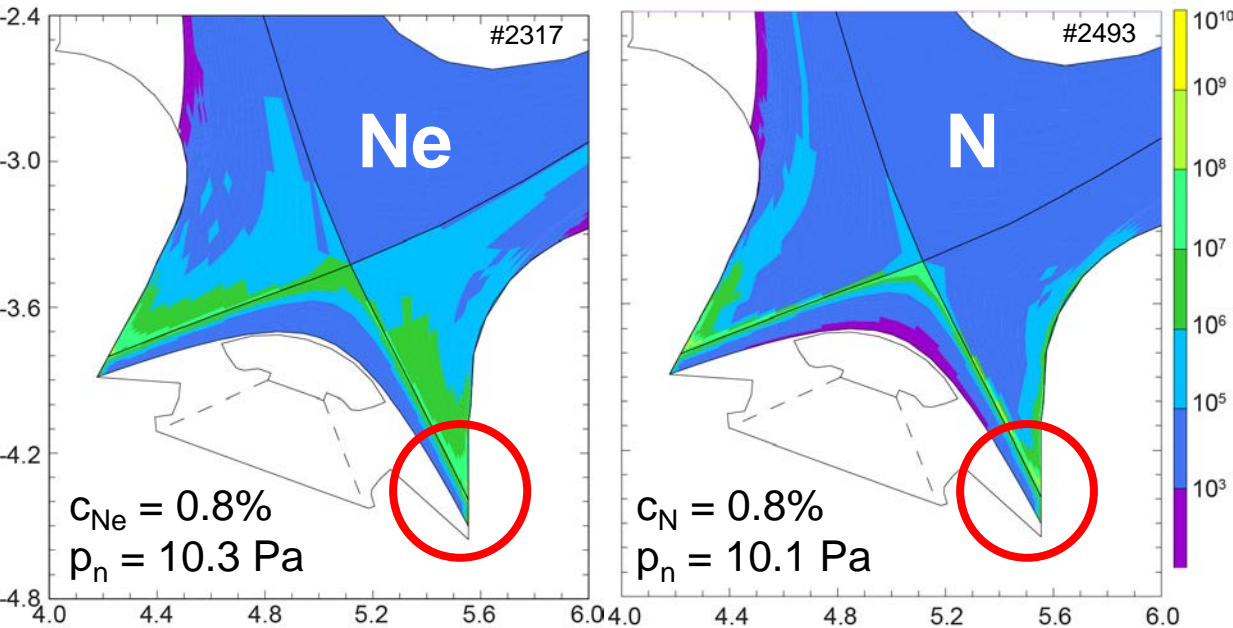
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Divertor radiation distribution

Total radiated power (Wm^{-3})



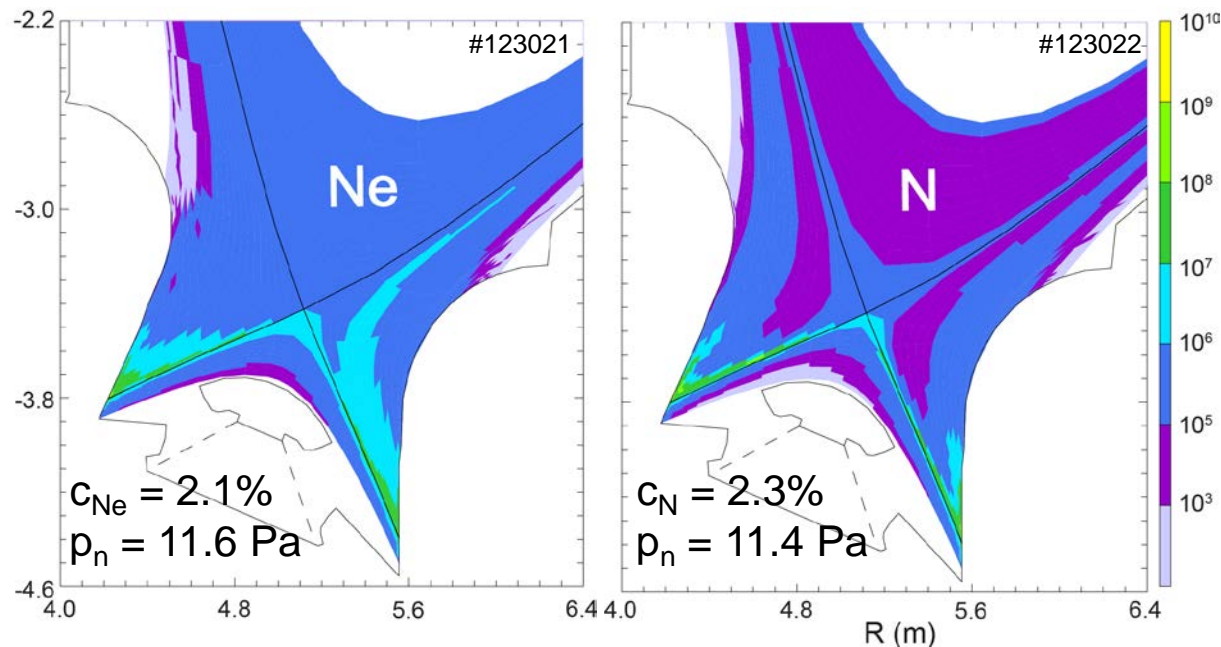
$P_{\text{RAD,DIV}} = 56.6 \text{ MW}$
41.3 (Ne) + 15.3 (D)

$P_{\text{RAD,DIV}} = 54.0 \text{ MW}$
38.6 (N) + 15.4 (D)

- Ne radiation more extended than N
 - Expected from differences in ionization potential
 - But still mostly confined to divertor volume
- Compression:
 - $n_{\text{Z,osp}}/n_{\text{Z,omp}}$
 - ~ 100 (N), ~ 30 (Ne)

Divertor radiation distribution: with drifts

Total radiated power (Wm^{-3})



$P_{\text{RAD,DIV}} = 53.4 \text{ MW}$
38.7 (Ne) + 14.2 (D)

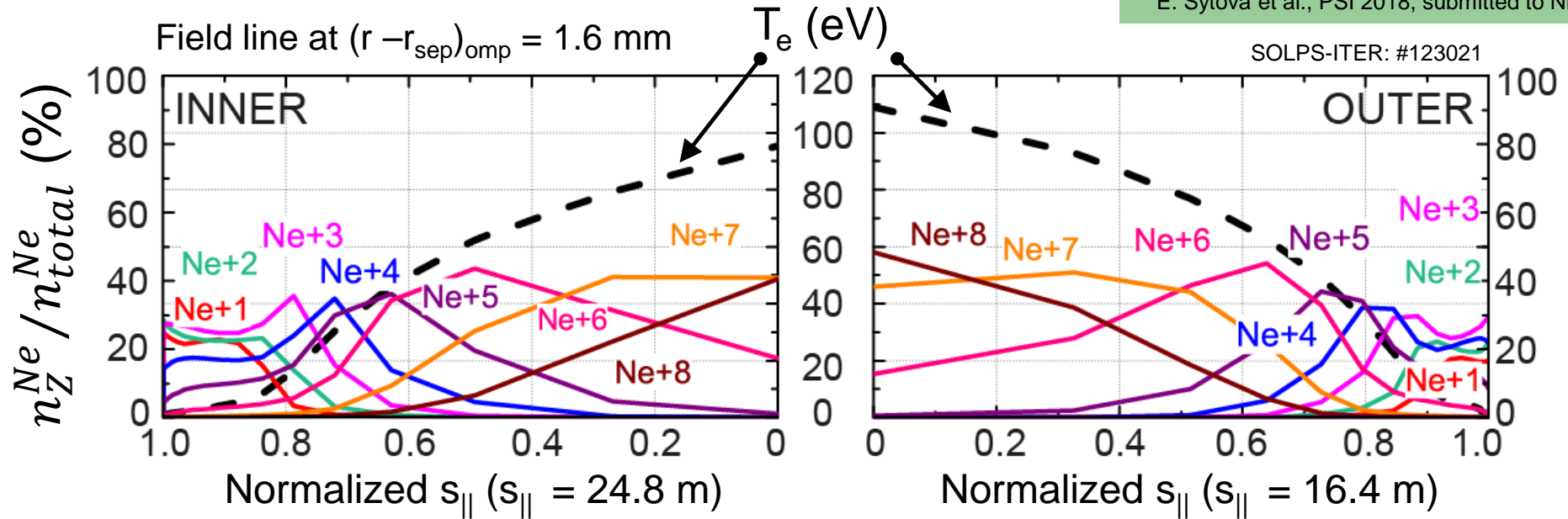
$P_{\text{RAD,DIV}} = 54.0 \text{ MW}$
44.7 (Ne) + 9.2 (D)

- New results from SOLPS-ITER
 - $P_{\text{IN}} = 100 \text{ MW}$
 - Matched Ne, N cases
 - H-mode pedestal
- Similar to SOLPS-4.3
 - Drift effects not important at high p_n
 - High magnetic field and better “divertor impurity screening”

E. Sytova et al., submitted to NME

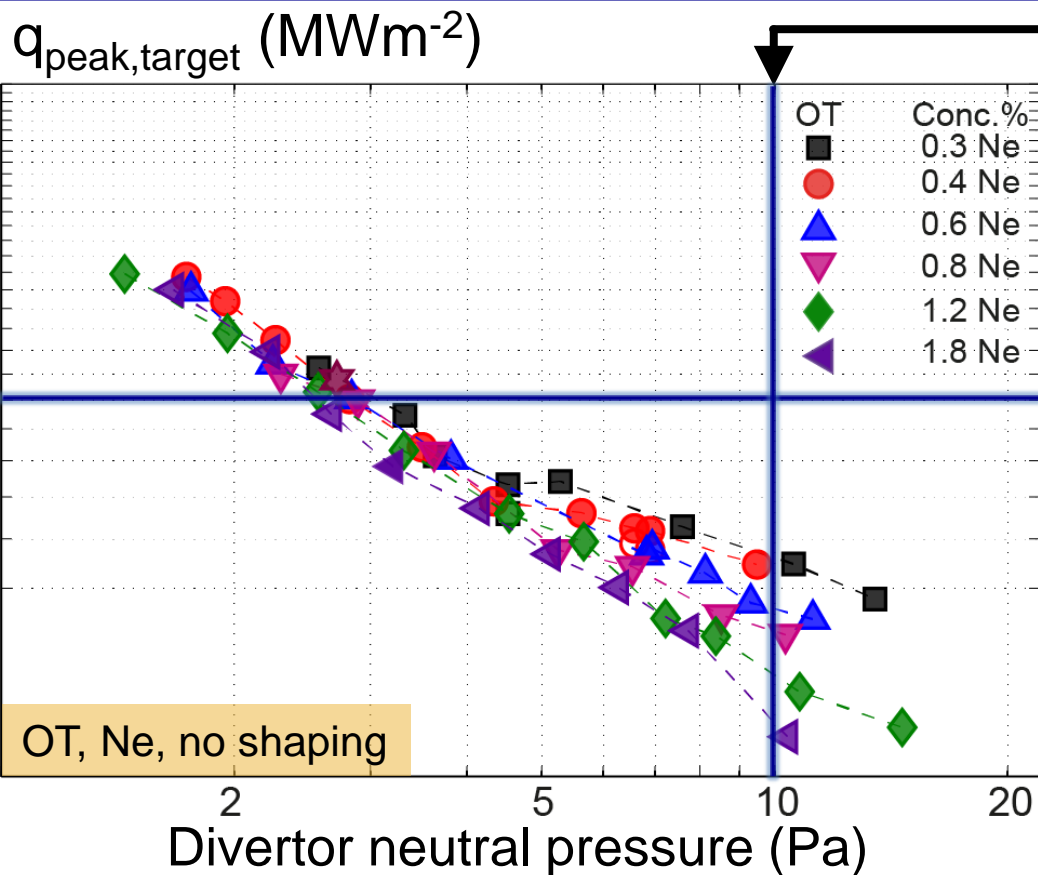
Impurity charge state distribution

E. Sytova et al., PSI 2018, submitted to NME



- $\sim 87\%$ of the divertor radiation comes from $\text{Ne}^{+3} \rightarrow \text{Ne}^{+6}$
 - Well confined in the divertor region $\rightarrow T_e$ high enough, far enough
 - Ne fully stripped in pedestal region and cannot radiate

Operating window in peak power flux density

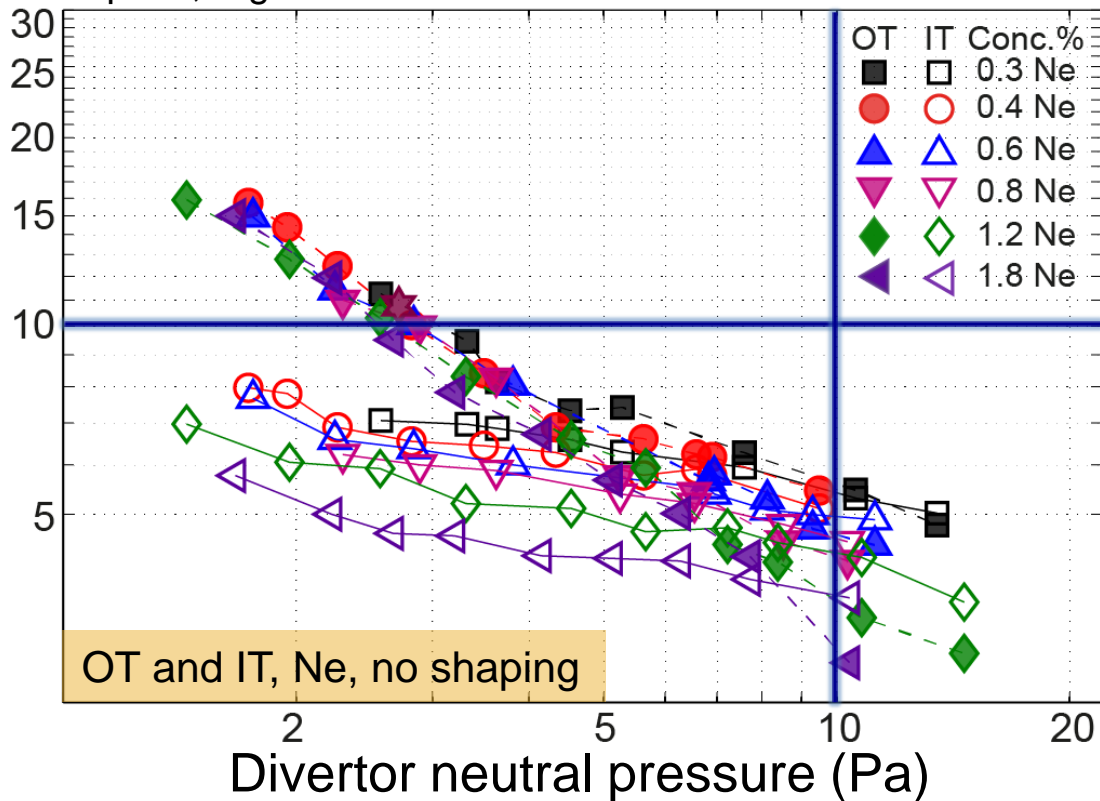


Approximate
“Detachment limit”
(see later)

“Historical” power
handling limit – will be
higher in reality

Operating window in peak power flux density

$q_{\text{peak,target}}$ (MWm^{-2})

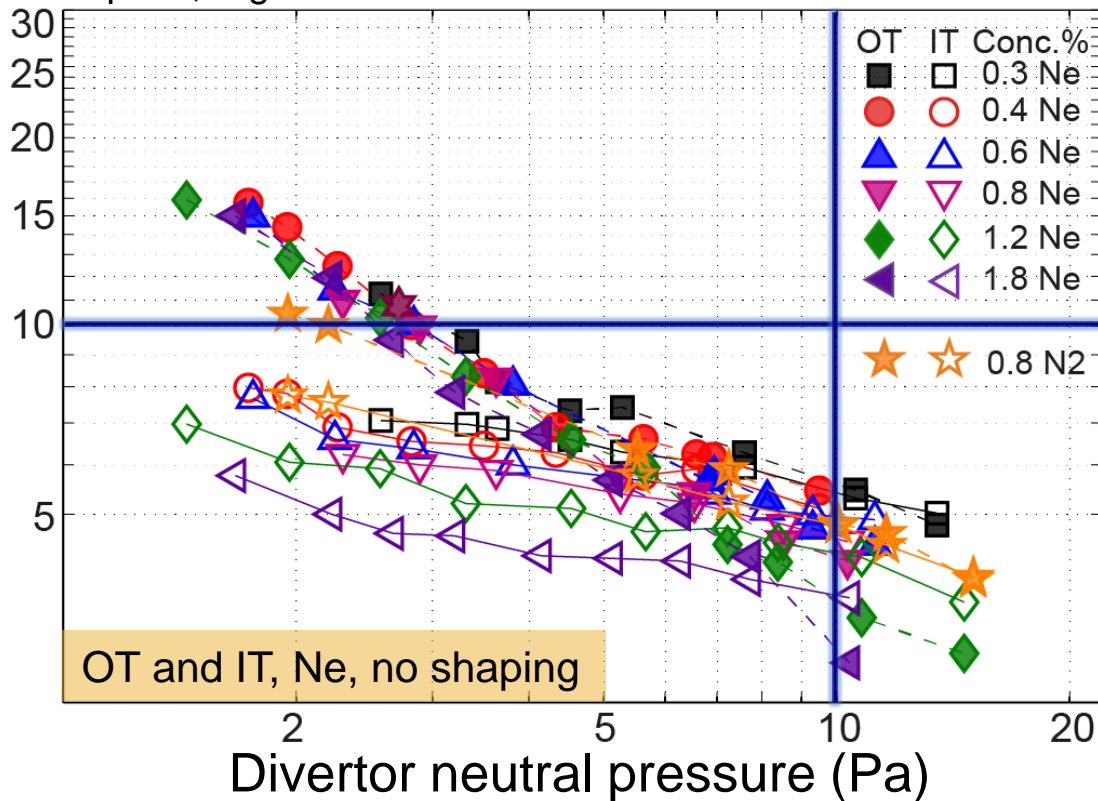


- Out-in asymmetry reduces at high p_{neut}
 - Strong neutral convection from inner to outer through private flux region balances ion flow from outer to inner target through the SOL → not a drift effect

A. A. Pshenov et al, PoP 24 (2018) 072508

Operating window in peak power flux density

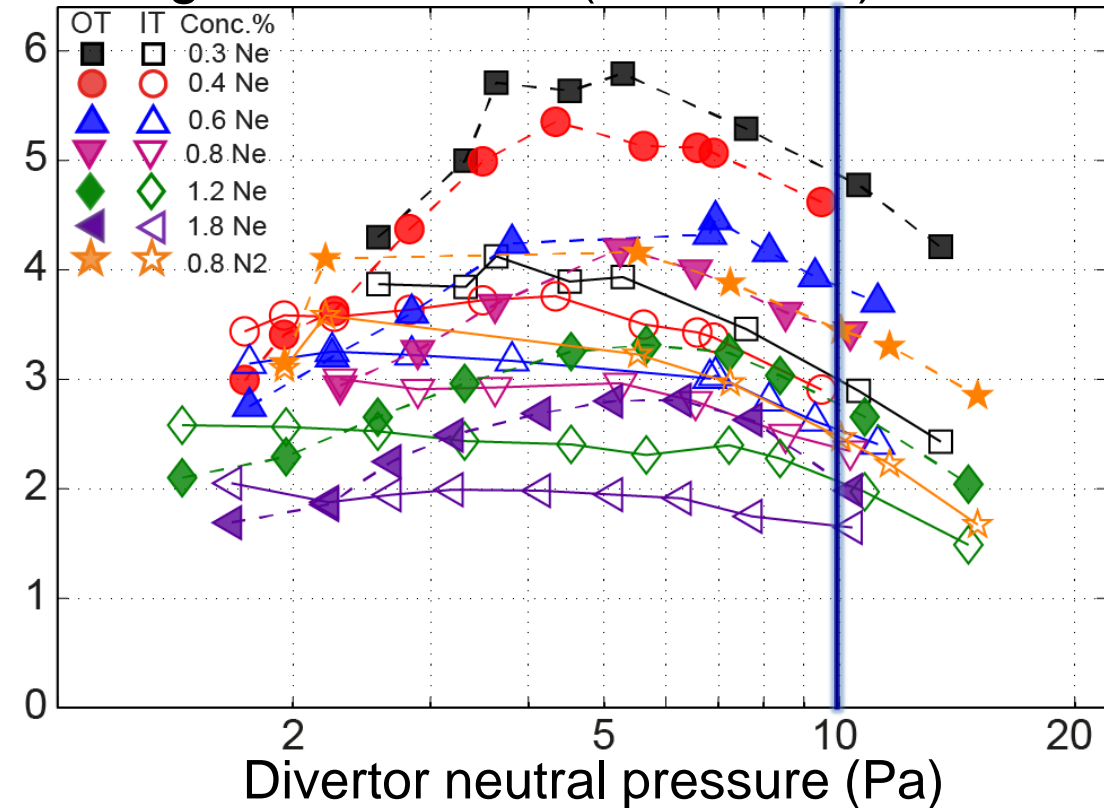
$q_{\text{peak,target}}$ (MWm^{-2})



- Nitrogen points overlay well with Ne cases
 - Database more restricted for N but trends similar
 - Need 3-5x as much N than Ne in the code for given D fueling to obtain similar midplane impurity separatrix concentration

Integrated target ion fluxes

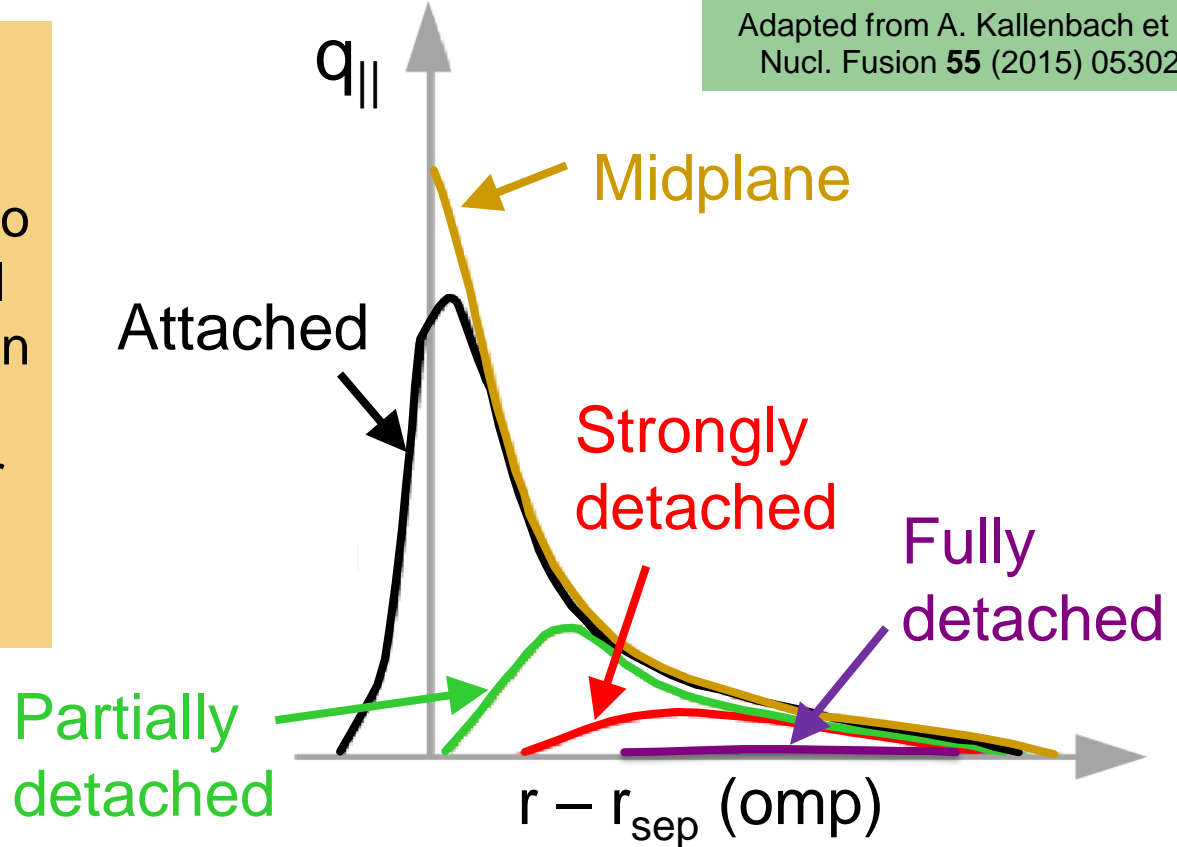
Integral ion current ($\times 10^{24}$ Ds $^{-1}$)



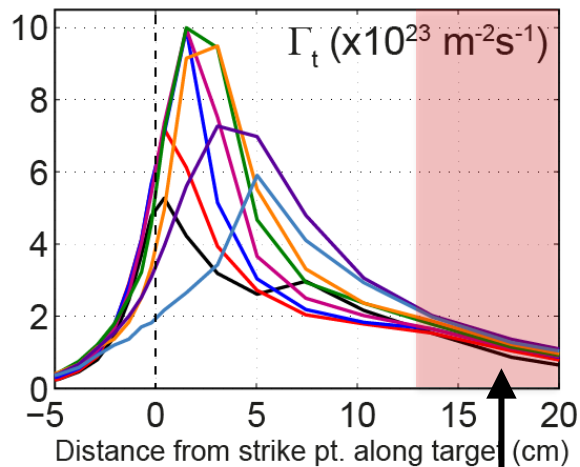
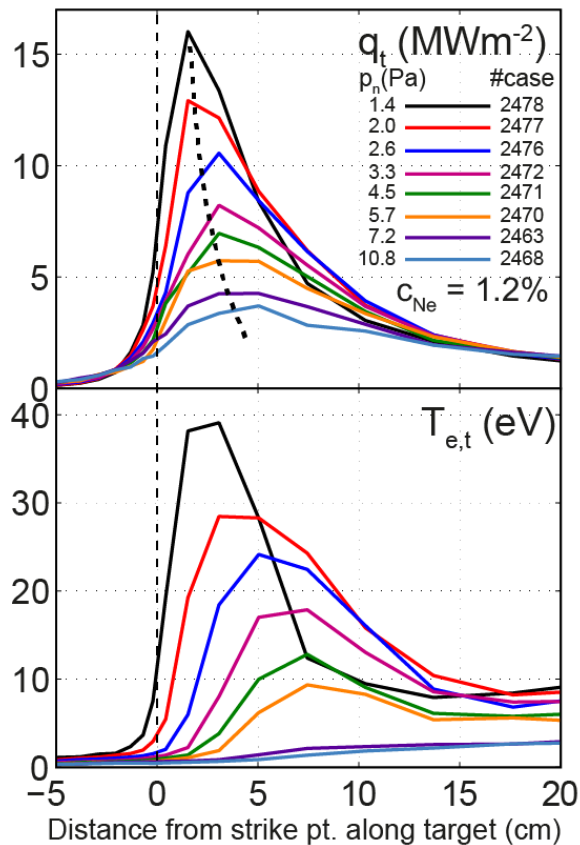
- Turnover in total plate current generally rather gentle
 - Loose criterion for “tolerable detachment” fixed as point at which integral flux reaches ~80% of peak value after rollover (based historically on discussions with JET) → happens typically near $p_n \sim 10$ Pa

Classic picture of detachment

This is where ITER expects to operate (bring down the ion flux to the plate since potential released in recombination at the target is a major contributor to the power load – see talk by D. Reiter)



Detachment evolution: outer target example

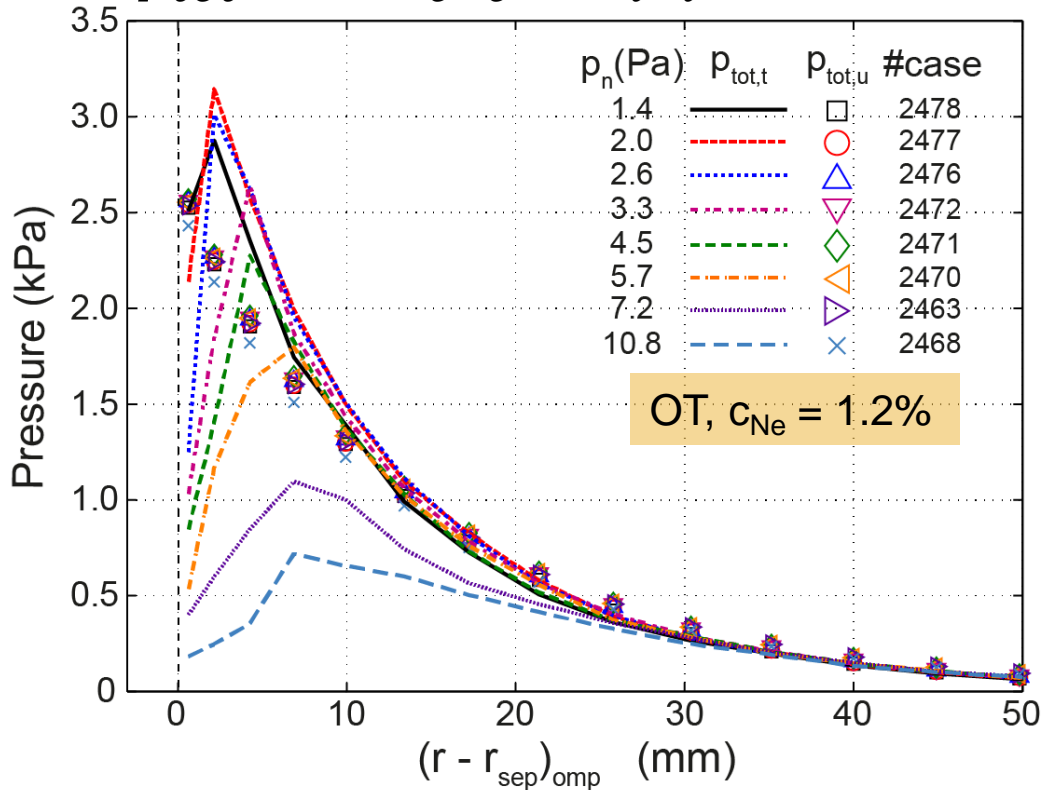


- “Classic” evolution from high recycling to partially detached state
 - He pumping improves with increased p_n

Avoid “complete” detachment \rightarrow keep finite ion flux in outer part of the SOL to maintain sufficient neutral plugging

Total pressure-momentum losses

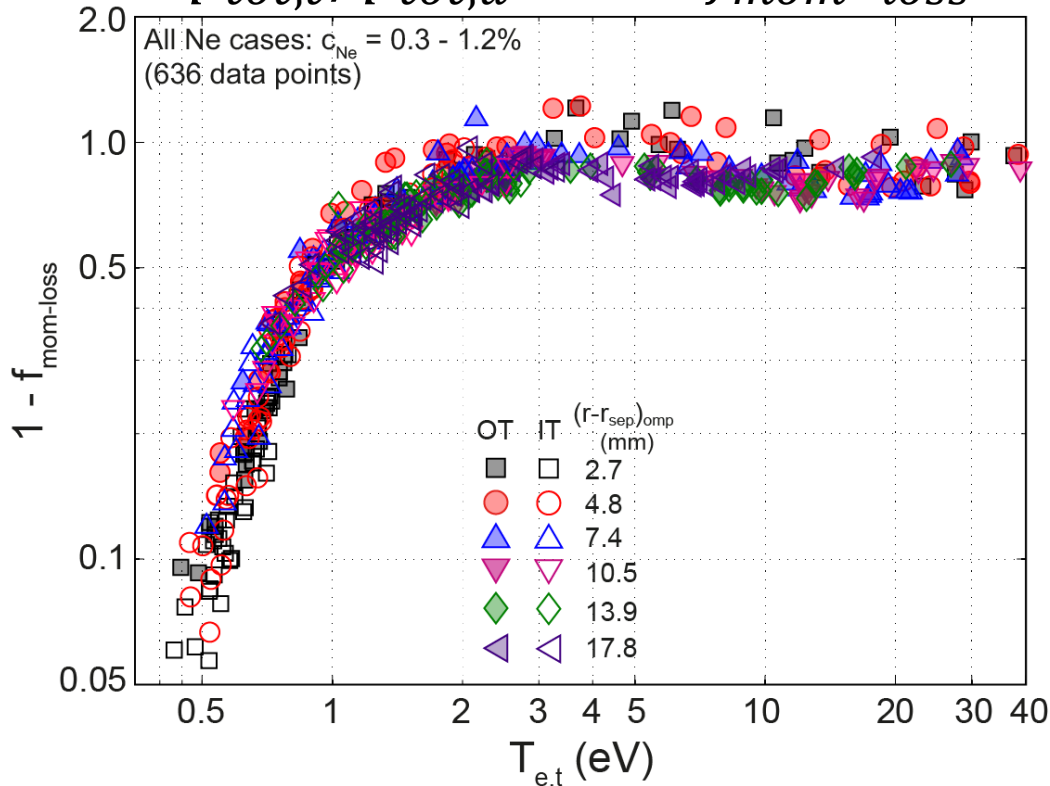
$$p_{tot} = k(n_e T_e + n_i T_i)(1 + M^2)$$



- Pressure loss downstream as p_n increases
 - Upstream p_{tot} unaffected by downstream conditions (as for λ_q)
 - Beyond region of pressure loss, upstream and downstream profiles overlap

Pressure-momentum losses vs. $T_{e,t}$

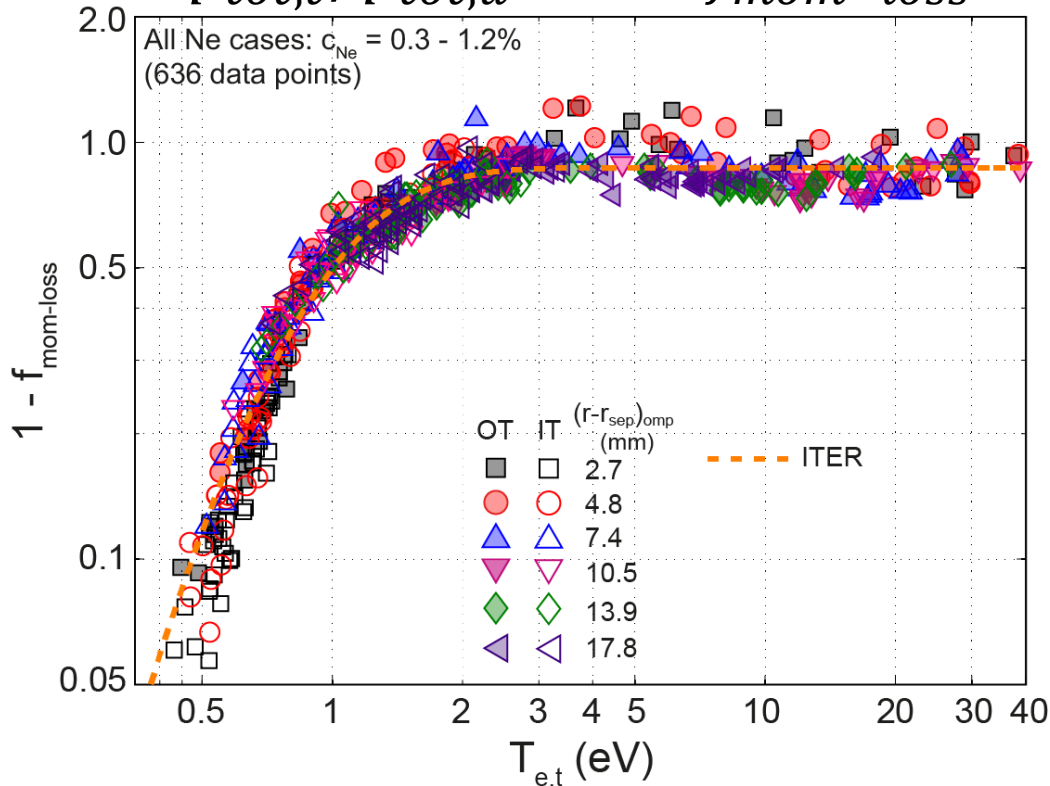
$$p_{tot,t}/p_{tot,u} = 1 - f_{mom-loss}$$



- Momentum loss strong function of $T_{e,t}$
 - Same at both IT and OT
 - Loss starts at $T_{e,t} \sim 3$ eV
 - Strong below $T_{e,t} \sim 1$ eV
 - Implies strong role for volume recombination → important for ITER (magnetic field incidence angles are high)

Pressure-momentum losses vs. $T_{e,t}$

$$p_{tot,t}/p_{tot,u} = 1 - f_{mom-loss}$$



- Momentum loss strong function of $T_{e,t}$
 - ITER SOLPS-4.3 simulation database well fitted by functional form proposed by Stangeby*:

$$1 - f_{mom}^{tot} = A(1 - e^{-T_{et}/T^*})^n$$

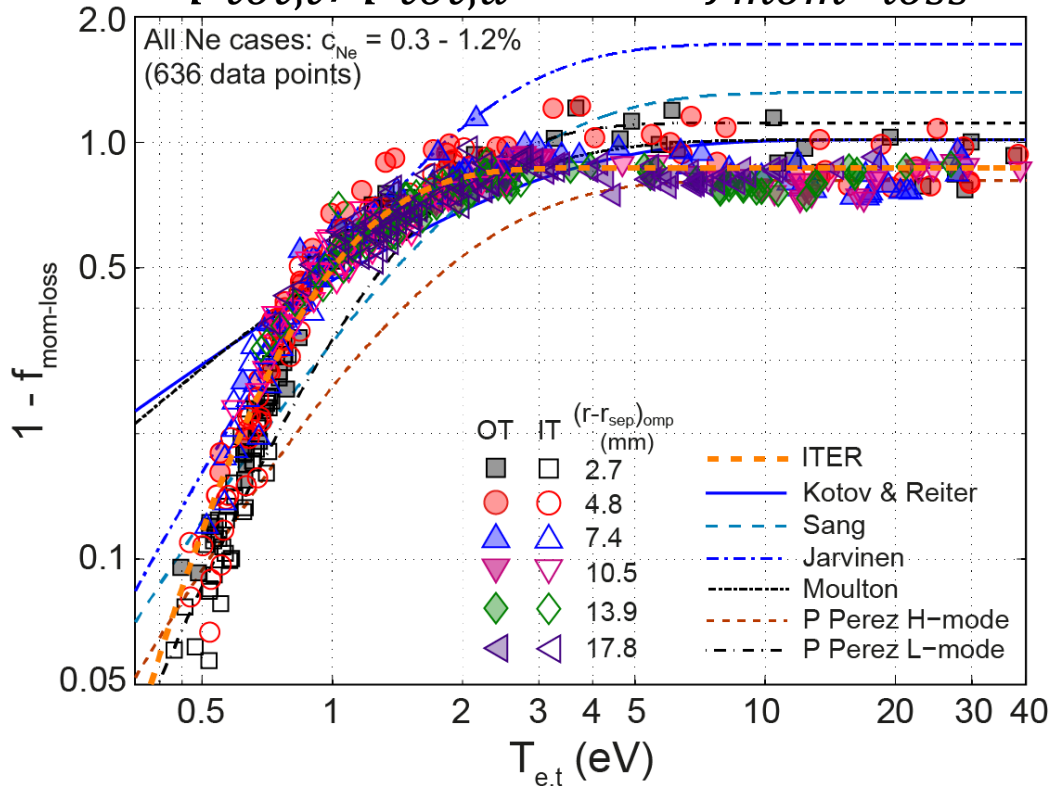
$$A = 0.87, T^* = 0.43, n = 5.30$$

$$R^2 = 0.96$$

*P. C. Stangeby, PPCF **60** (2018) 044022

Pressure-momentum losses vs. $T_{e,t}$

$$p_{tot,t}/p_{tot,u} = 1 - f_{mom-loss}$$

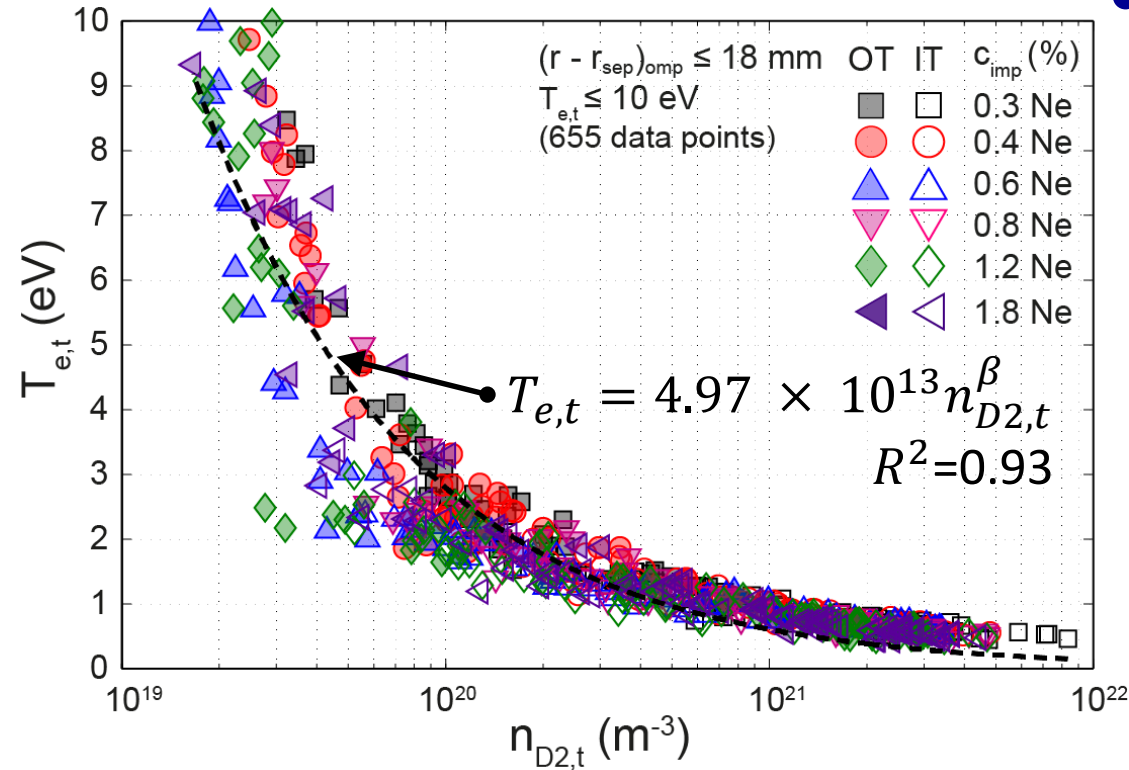


- Generally steeper trend than similar fits to other code studies*

- BUT note: “balance analysis” now underway indicates that particle removal by VR not necessarily the dominant process
- Neutral atoms and molecules created by VR responsible for up to half of $f_{mom-loss}$

*P. C. Stangeby, PPCF **60** (2018) 044022

Plate molecular density and $T_{e,t}$



- $T_{e,t}$ and $n_{D2,t}$ tightly correlated ($T_{e,t} \lesssim 10 \text{ eV}$)
 - Similar to findings from other code studies
 - $n_{D2,t}$ related to flux amplification at targets and hence to $f_{\text{mom-loss}} \rightarrow$ expect link between $n_{D2,t}$ and $T_{e,t}$
 - Currently investigating link between $T_{e,t}$ and volumetric power losses

V. Kotov, D. Reiter, PPCF **51** (2009) 115002, P. C. Stangeby, PPCF **60** (2018) 044022

Factors influencing peak power density

λ_q narrower than we have assumed? \rightarrow best current exptl. scaling gives $\lambda_q \sim 1$ mm for ITER at $I_p = 15$ MA



Neoclassical ion drift model very good match to scaling

R. J. Goldston, NF 52 (2012) 013009

No inconsistency with SOLPS-ITER with drifts and reduced transport

V. Rozhansky et al., PPCF 60 (2018) 035001

Also holds at ITER values of B_{pol} on C-Mod

D. Brunner et al., NF 58 (2018) 094002

PFC shaping

Fluid drifts? (SOLPS-4.3 has no drift capability)

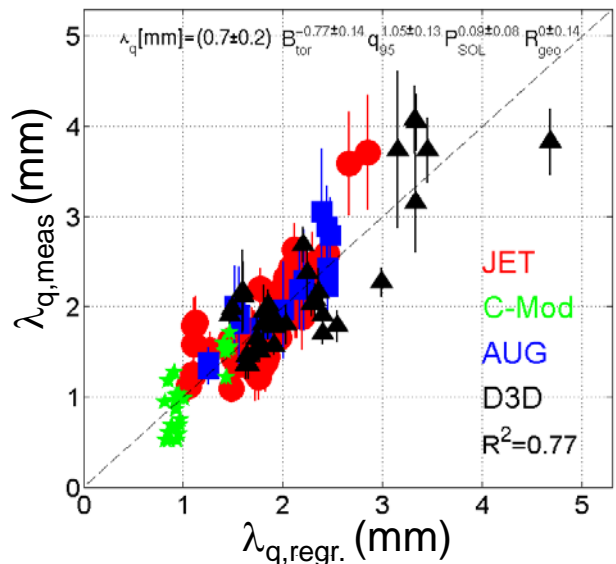
Insufficient numerical simulation grid resolution?



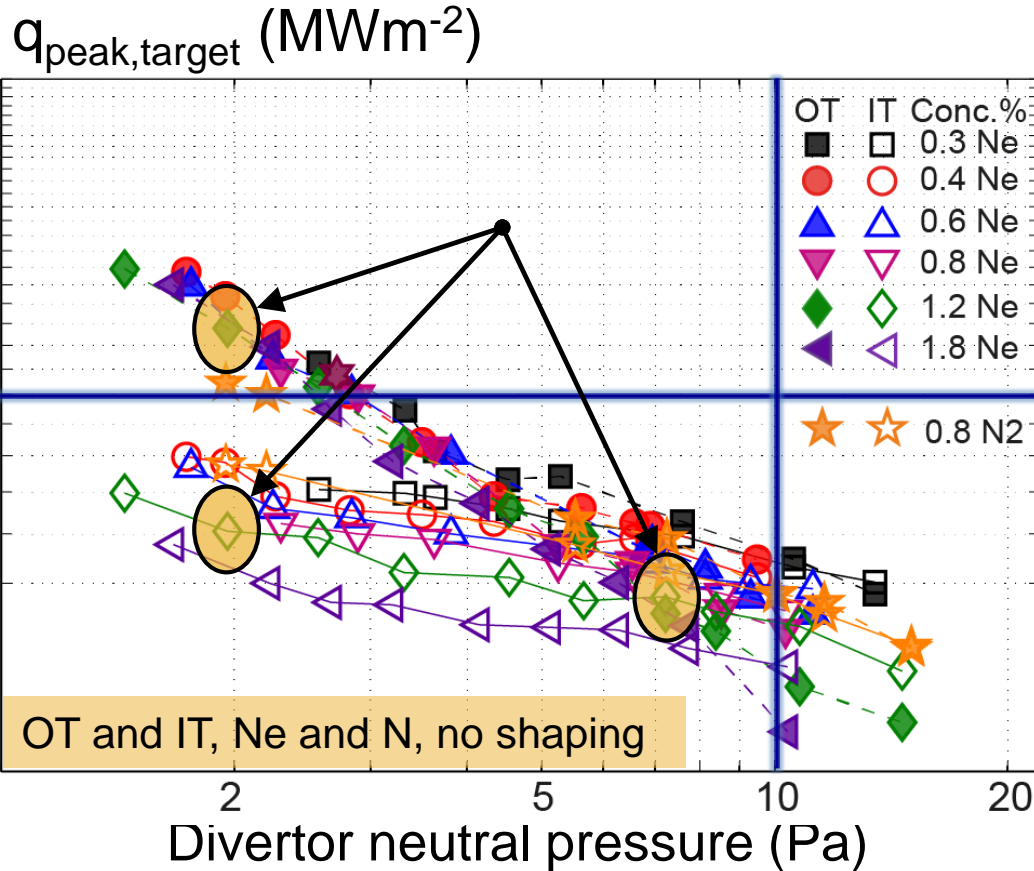
Being actively pursued at Univ. Leuven

K. Ghoos et al, submitted to NF

T. Eich, NF 53 (2013) 093031



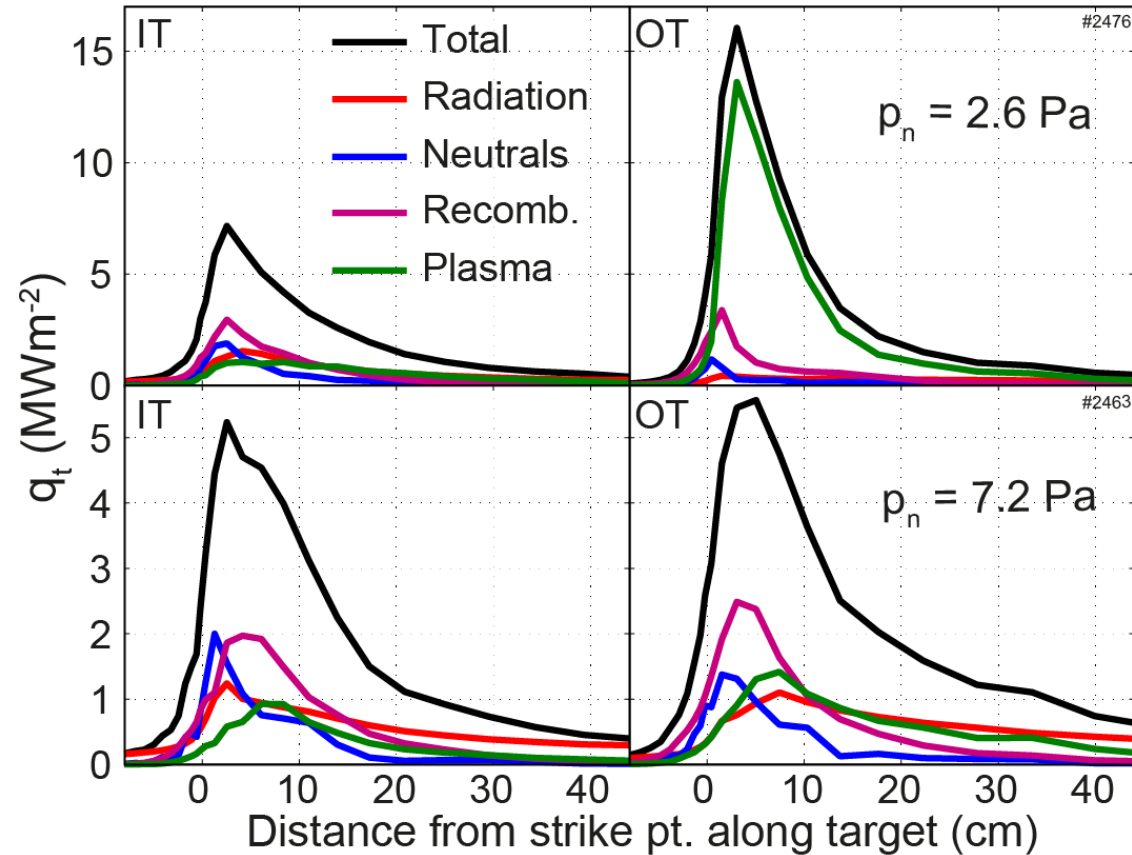
Impact of shaping



- Need to apply angle corrections for global target tilting and monoblock toroidal shaping **only** to thermal plasma components
 - Kinetic plasma plus potential energy of recombination at the plate:

$$\gamma n_{\text{et}} c_{\text{st}} T_{\text{et}} + n_{\text{et}} c_{\text{st}} E_{\text{pot}}$$

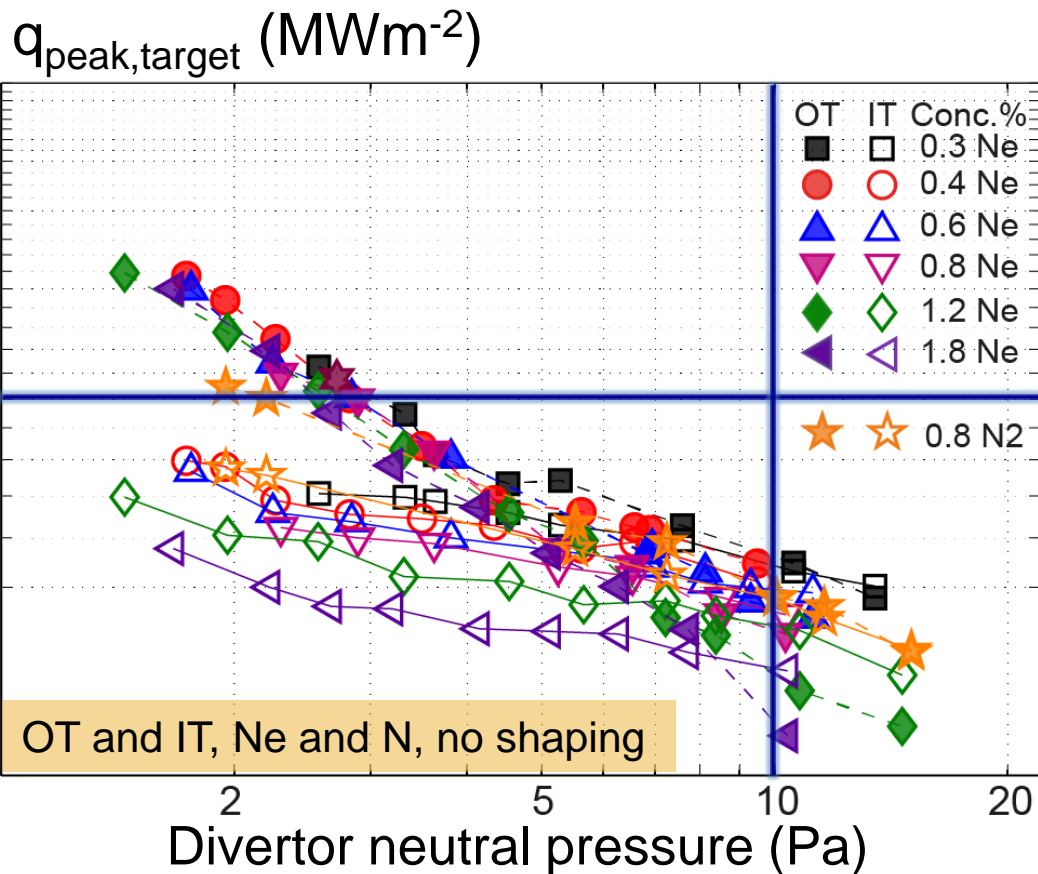
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- Need to apply angle corrections for global target tilting and monoblock toroidal shaping **only** to thermal plasma components
 - Kinetic plasma plus potential energy of recombination at the plate:

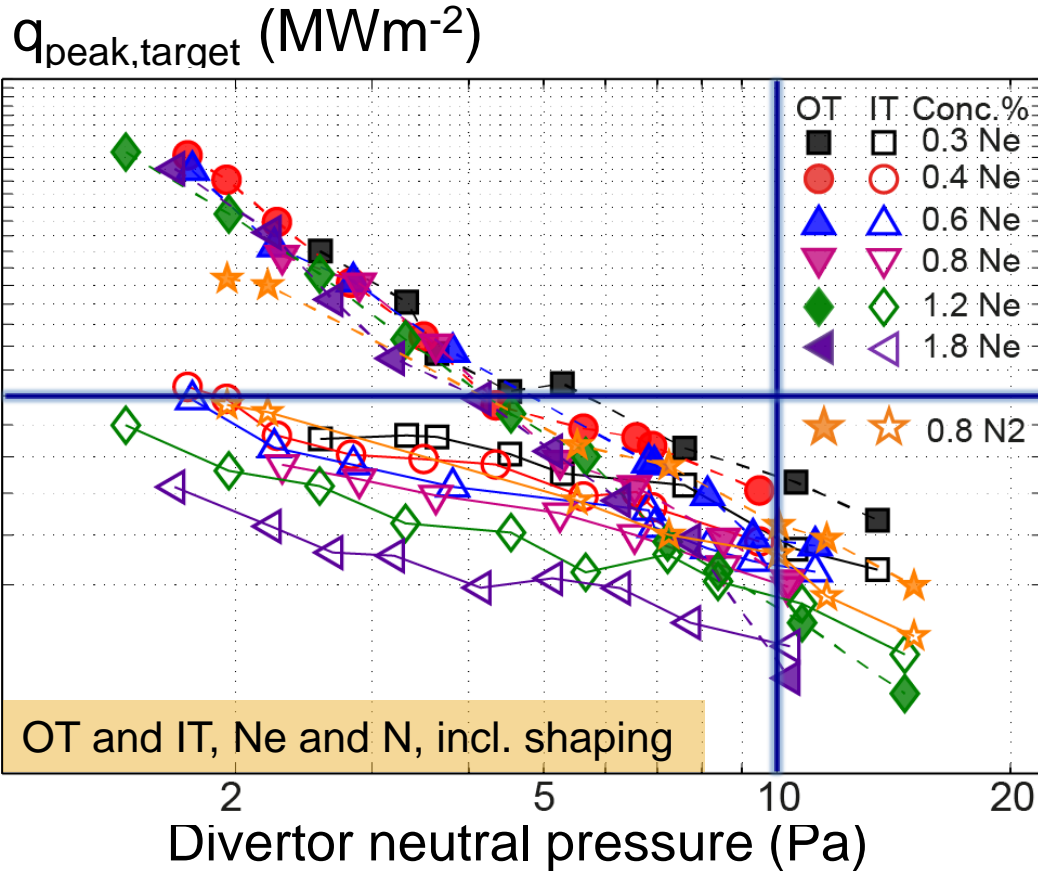
$$\gamma n_{et} c_{st} T_{et} + n_{et} c_{st} E_{pot}$$

Impact of shaping



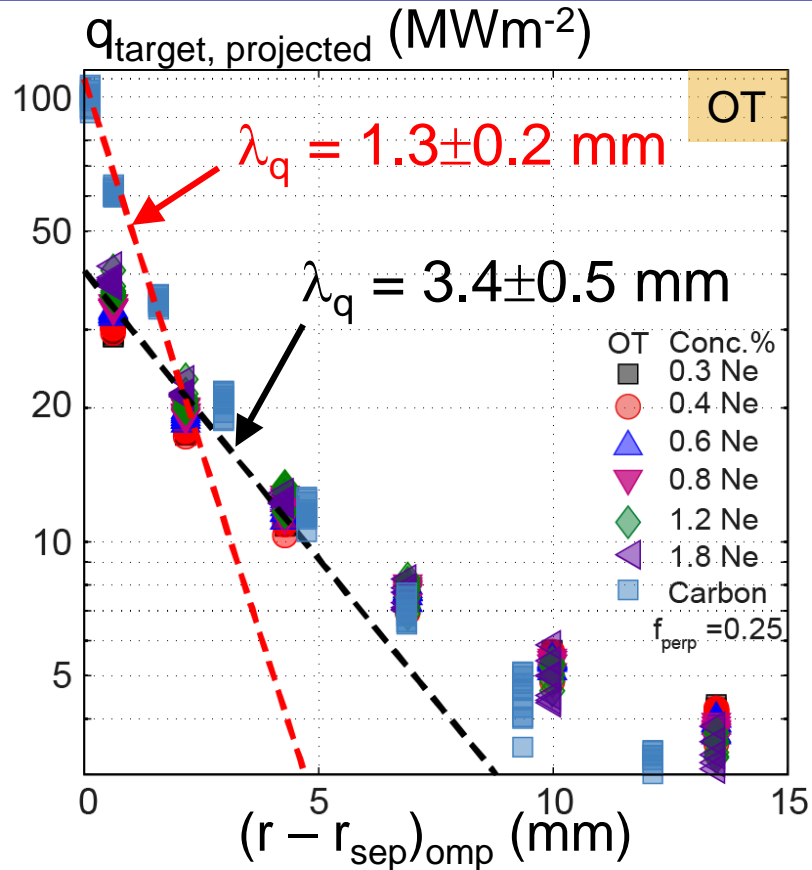
- Reminder, no shaping

Impact of shaping



- Now with shaping
 - Effects less marked at high p_n (where thermal plasma contributions lower)

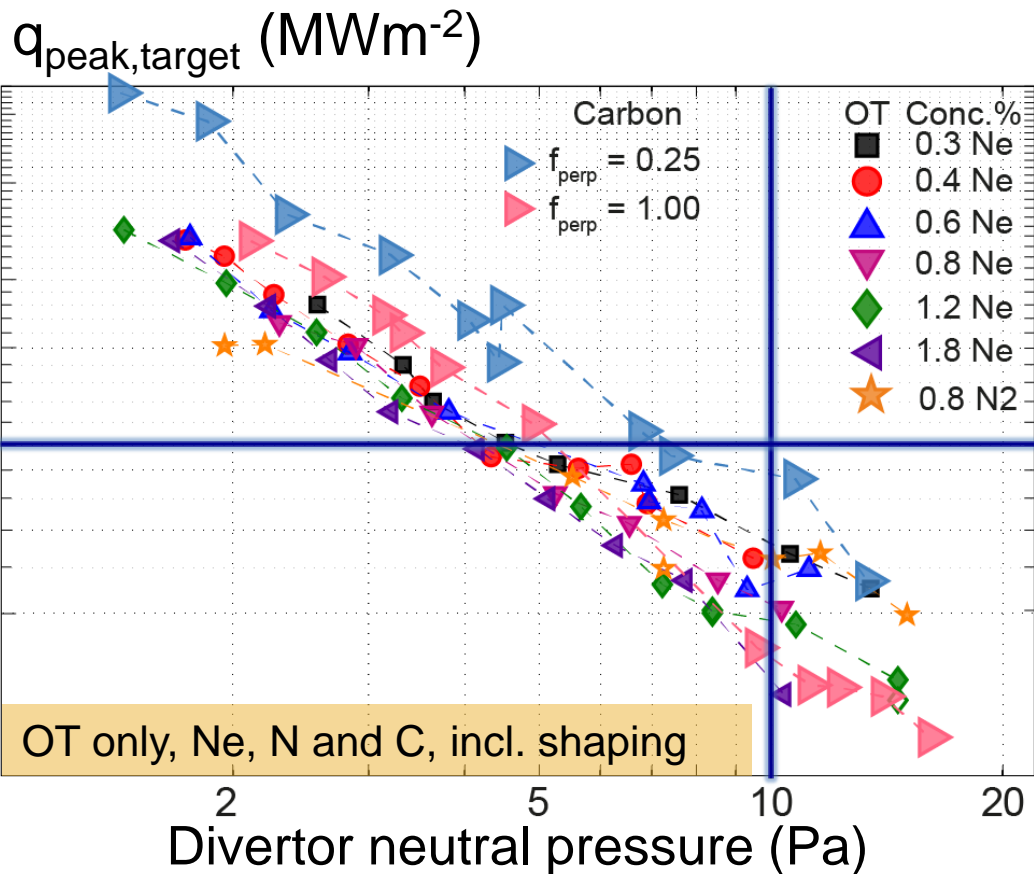
Impact of reduced transport



- Reduce D_{\perp} , χ_{\perp} by factor 4 compared to baseline:
 $D_{\perp} = 0.075 \text{ m}^2\text{s}^{-1}$
 $\chi_{\perp} = 0.25 \text{ m}^2\text{s}^{-1}$
 - Only old carbon divertor cases*
 - New SOLPS-ITER runs for Be/W underway
 - $\lambda_q = 1.3 \text{ mm}$ close to experimental scaling for ITER at 15 MA

*A. S. Kukushkin et al. JNM 438 (2013) S203

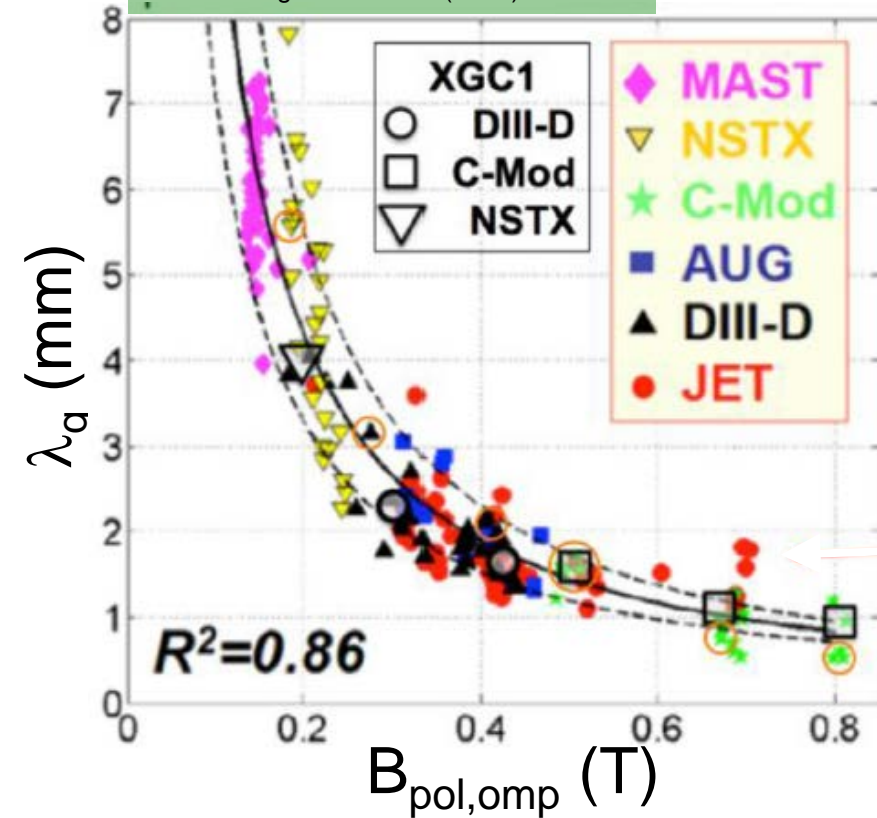
Impact of reduced transport and shaping



- Window for target heat flux density narrows to higher neutral pressure
 - Proximity to complete detachment threshold?
 - Upstream density limits and detachment stability?
 - R&D priority

What will be the true λ_q on ITER?

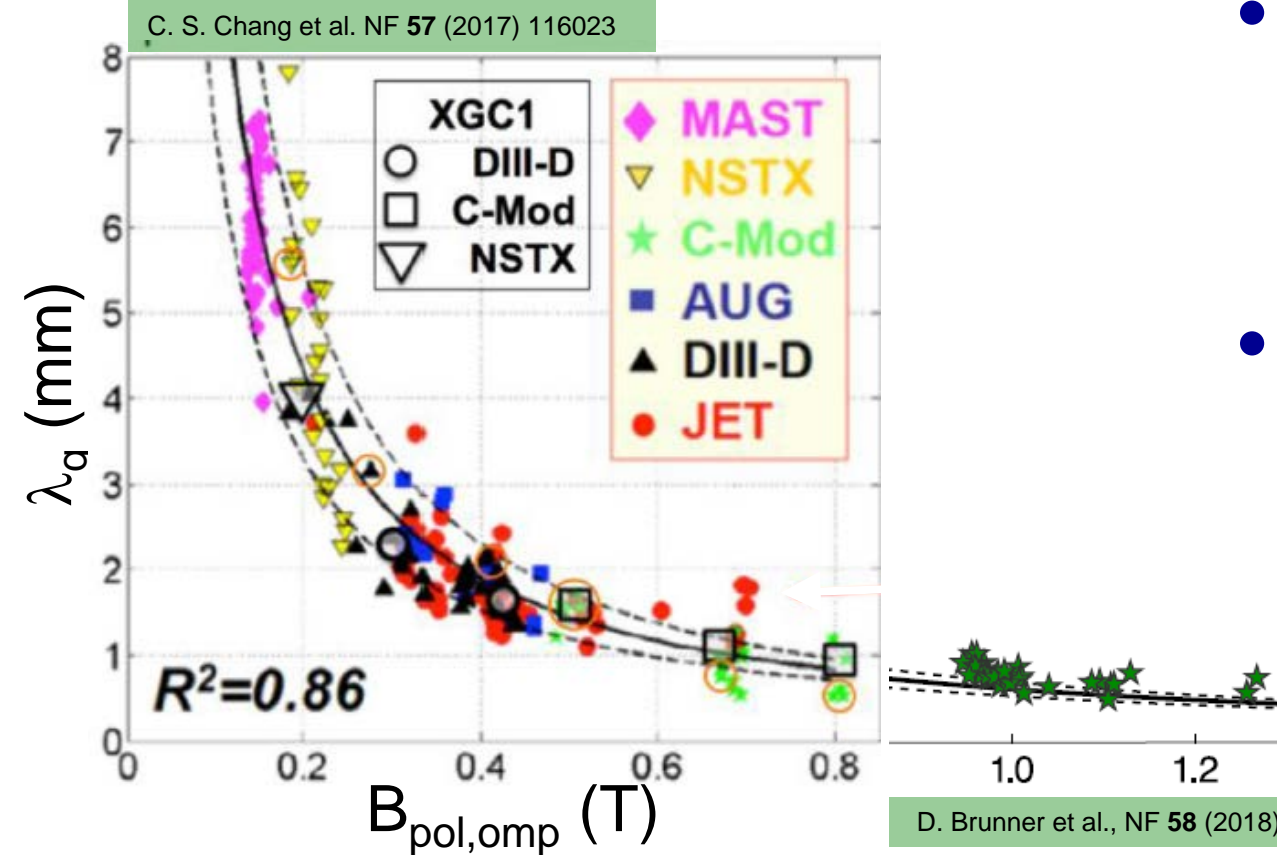
C. S. Chang et al. NF 57 (2017) 116023



- XGC1 electrostatic global gyrokinetic simulations match $\lambda_q \propto 1/B_{pol}$ scaling

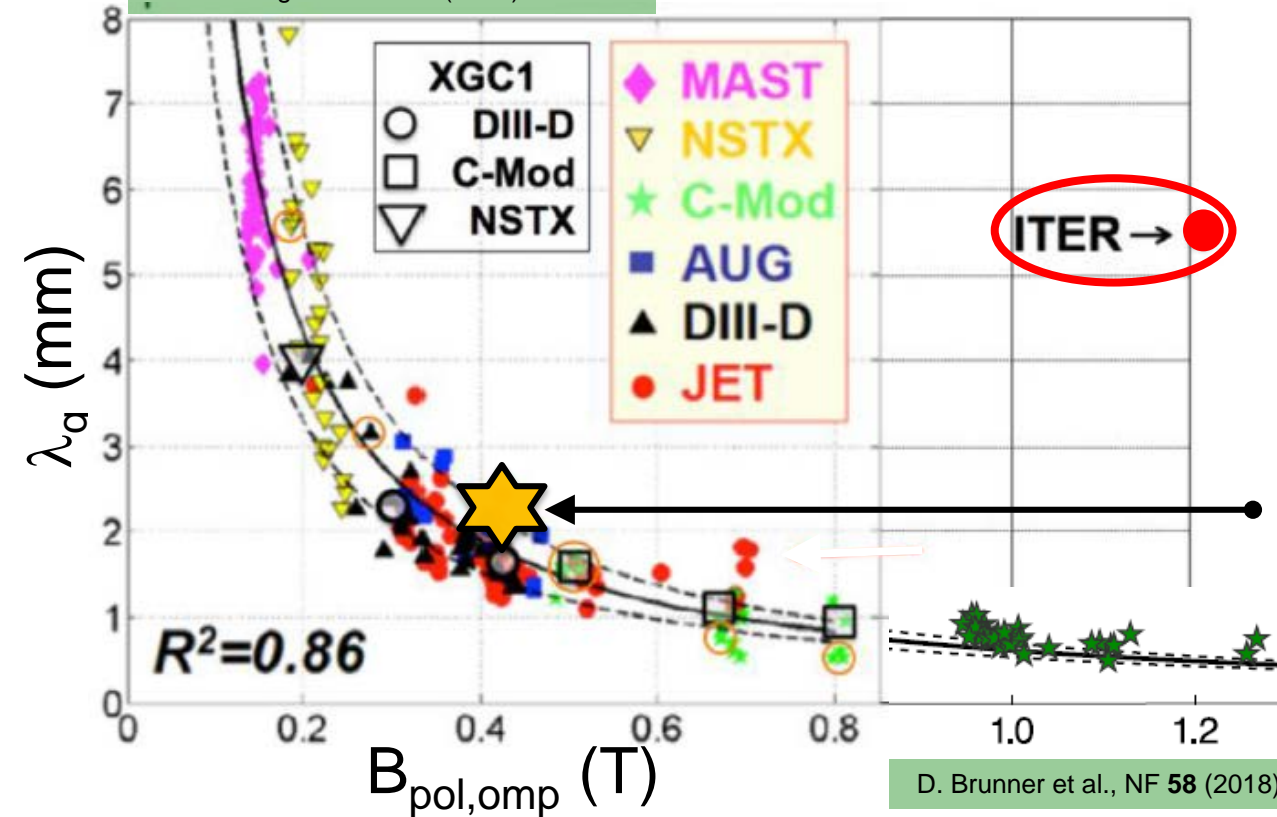
What will be the true λ_q on ITER?

- XGC1 electrostatic global gyrokinetic simulations match $\lambda_q \propto 1/B_{pol}$ scaling
- Data out to ITER B_{pol} on C-Mod continue to follow scaling



What will be the true λ_q on ITER?

C. S. Chang et al. NF 57 (2017) 116023

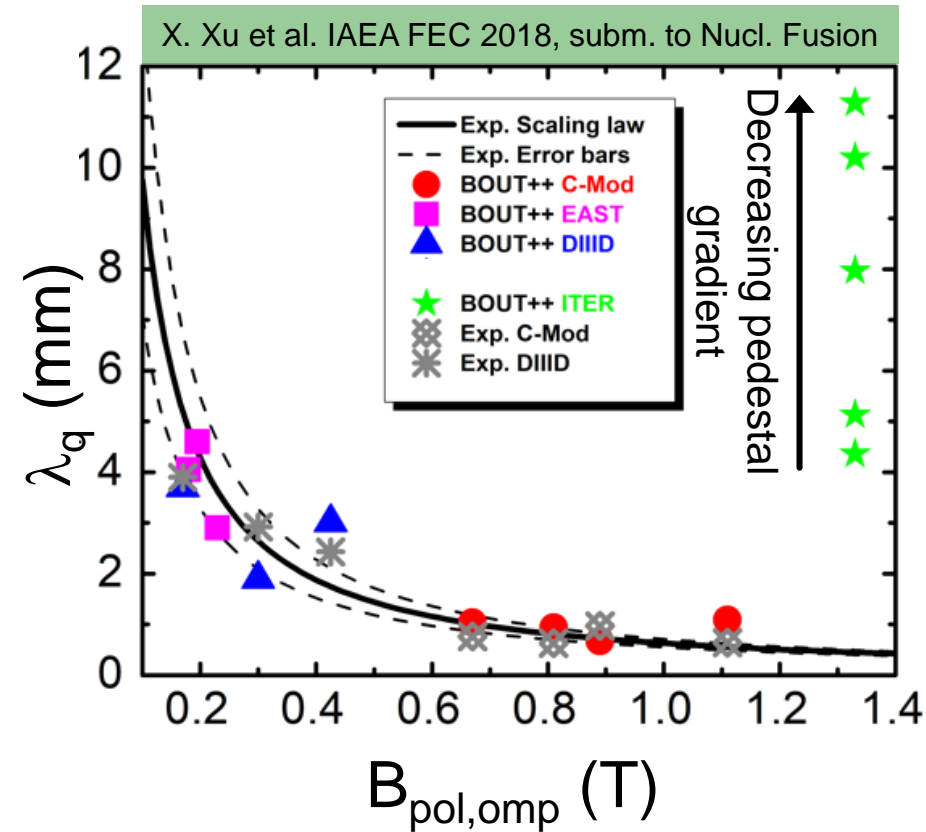


- But XGC1 dependence broken at ITER scale
 - Attributed to electron turbulence
 - Looks robust (for high I_p , B_T)
 - First XGC1 result for 5 MA ITER H-mode → follows empirical scaling

C. S. Chang et al., IAEA 2018

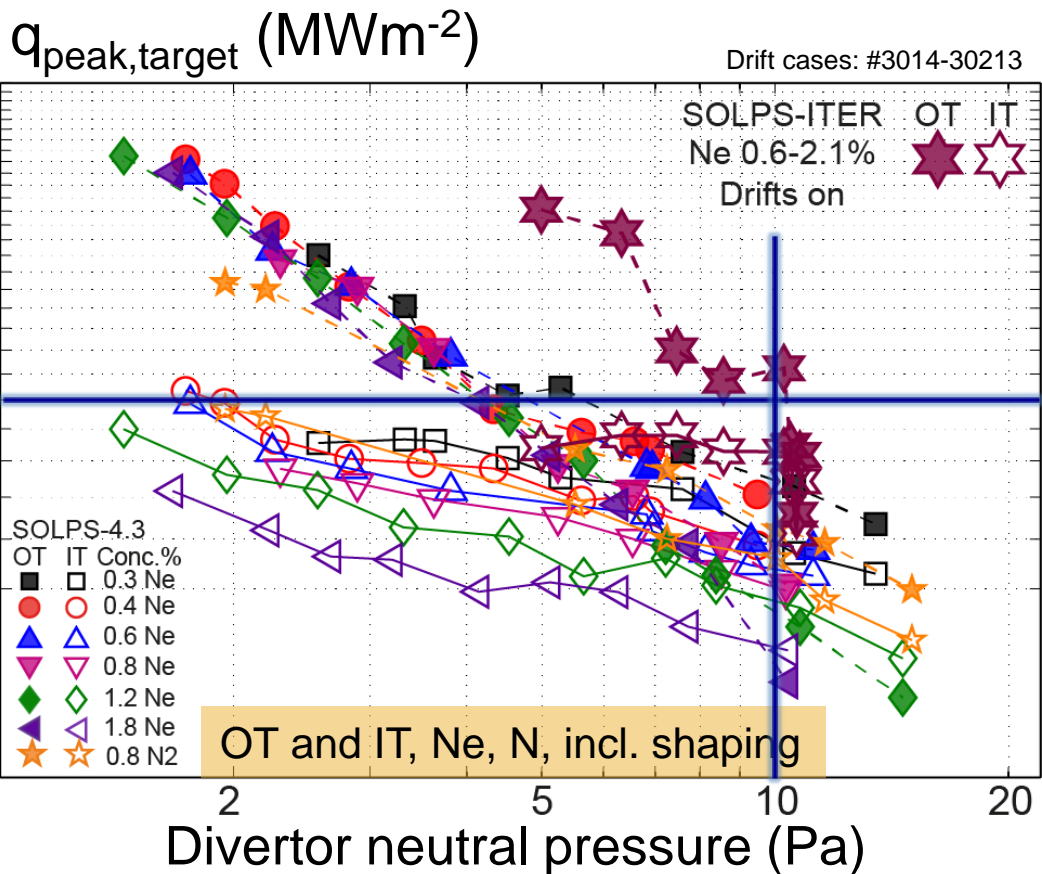
D. Brunner et al., NF 58 (2018) 094002

What will be the true λ_q on ITER?



- Very recent BOUT++ electromagnetic turbulence simulations find the same trends:
 - $\lambda_q \propto 1/B_{\text{pol}}$ for small devices
 - Broken at high I_p (B_{pol}) on ITER
 - Comparison of “transport BOUT” with “turbulence BOUT” shows a transition from magnetic drift to turbulence at given machine size
- This is a highly controversial and important issue for ITER and reactors

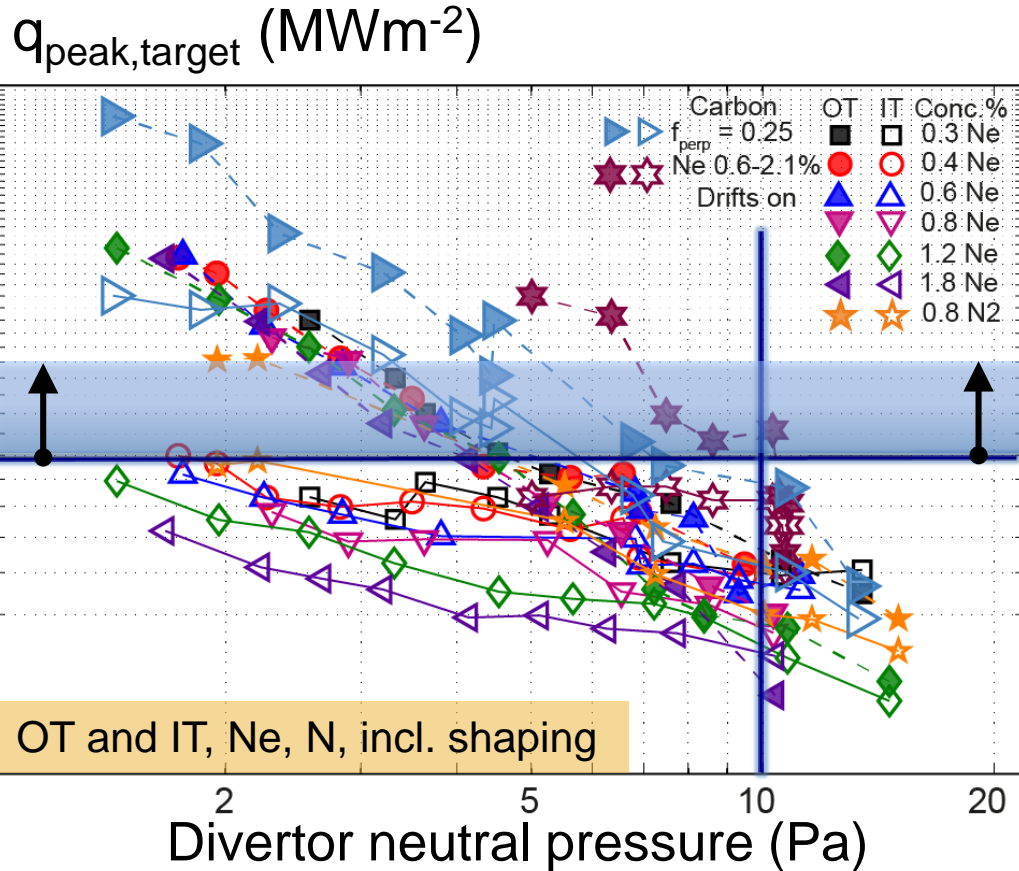
Impact of drifts



- Brand new results from “H-mode” SOLPS-ITER drift modelling*
- Strong impact on out-in asymmetry but effect reduced as detachment deepens
 - Neon very sensitive to drift-induced main ion flows
 - Only at the beginning of this study

*E. Kaveeva et al, in preparation for NF

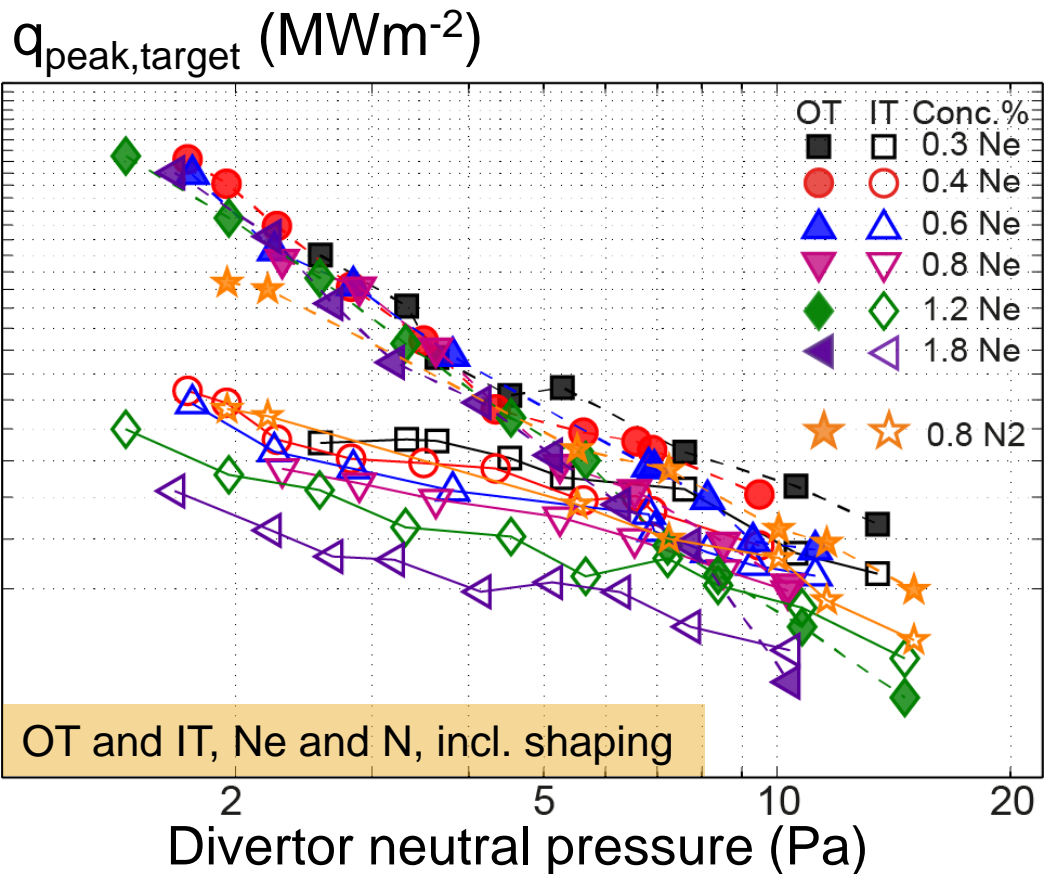
Being pushed into a corner?



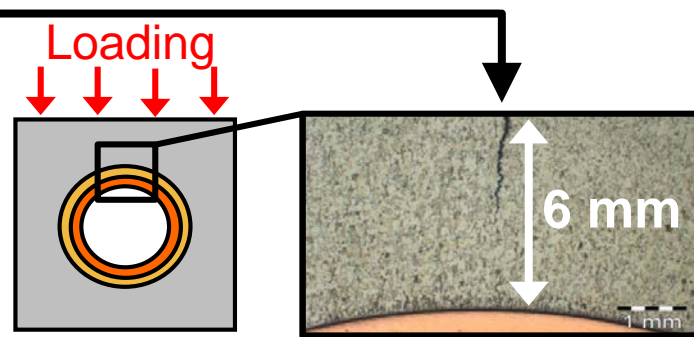
- Combined effect of all factors is push operation to higher divertor pressure
 - Good for He throughput
 - Potentially bad for detachment stability
 - New criterion for tolerable power handling helps a lot*
- Important to assess impact of operation at higher detachment degree

*G. De Temmerman et al, PPCF 60 (2018) 044018

What should be the true power load limit?

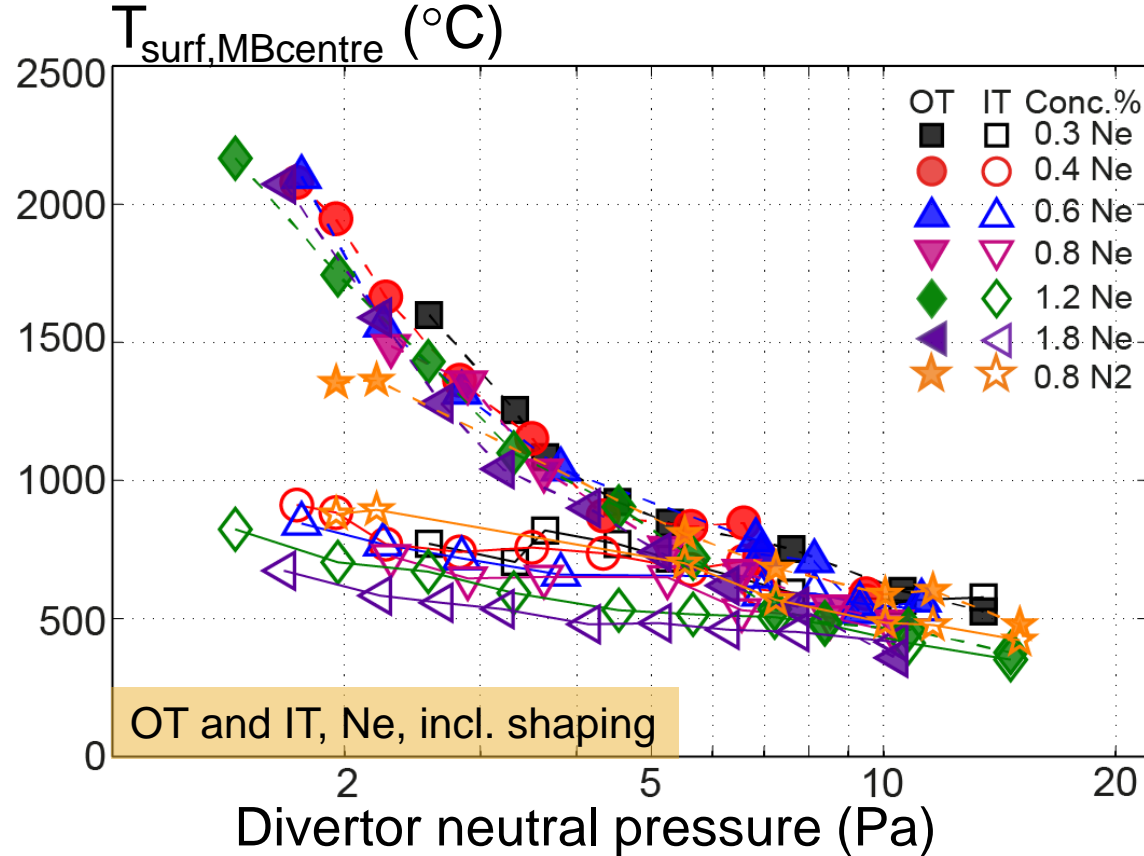


- Use FE simulations to transform q_{peak} to T_{surf}
 - Take value at monoblock centre → where cracking seen to start under high heat flux testing

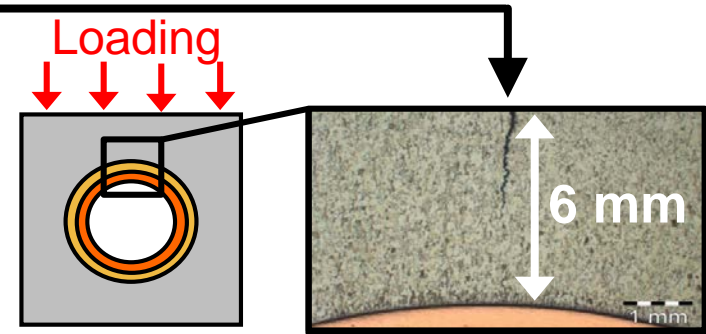


S. Panayotis et al., NME 12 (2017) 200

What should be the true power load limit?



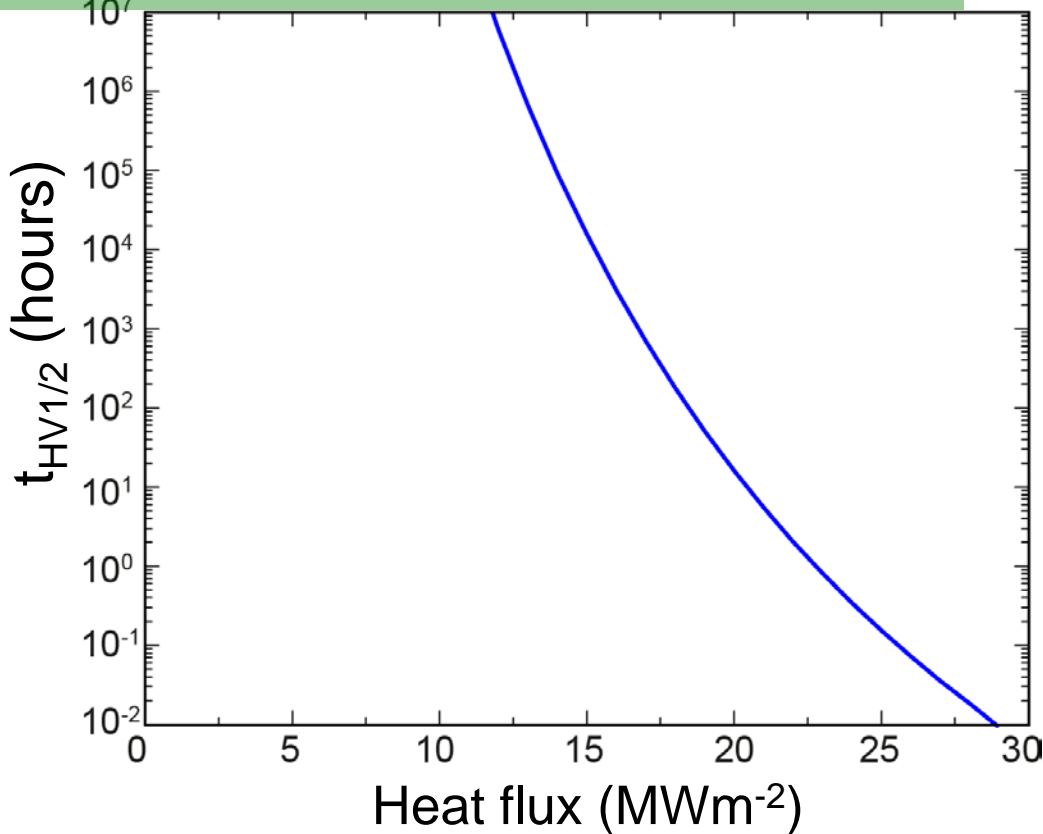
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S. Panayotis et al., NME 12 (2017) 200

What should be the true power load limit?

G. De Temmerman et al., PPCF **60** (2018) 044018, R. A. Pitts, NME submitted

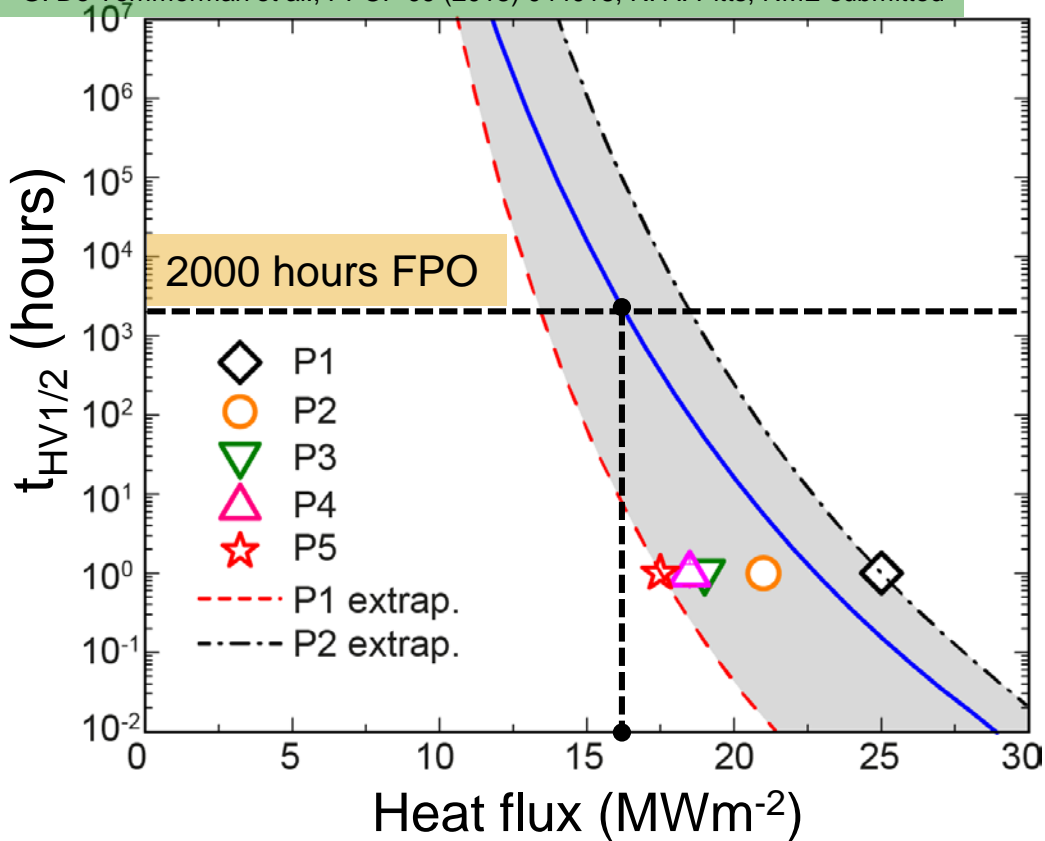


- “Operational budget” in $Q_{\text{peak,target}}$ defined by time required for hardness to drop by 50% at given depth below MB surface
 - ~2 mm recrystallization depth consistent with recent FEM modelling for crack onset due to low cycle fatigue¹

¹M. Li et al., Fus. Eng. Des. **101** (2018) 1

What should be the true power load limit?

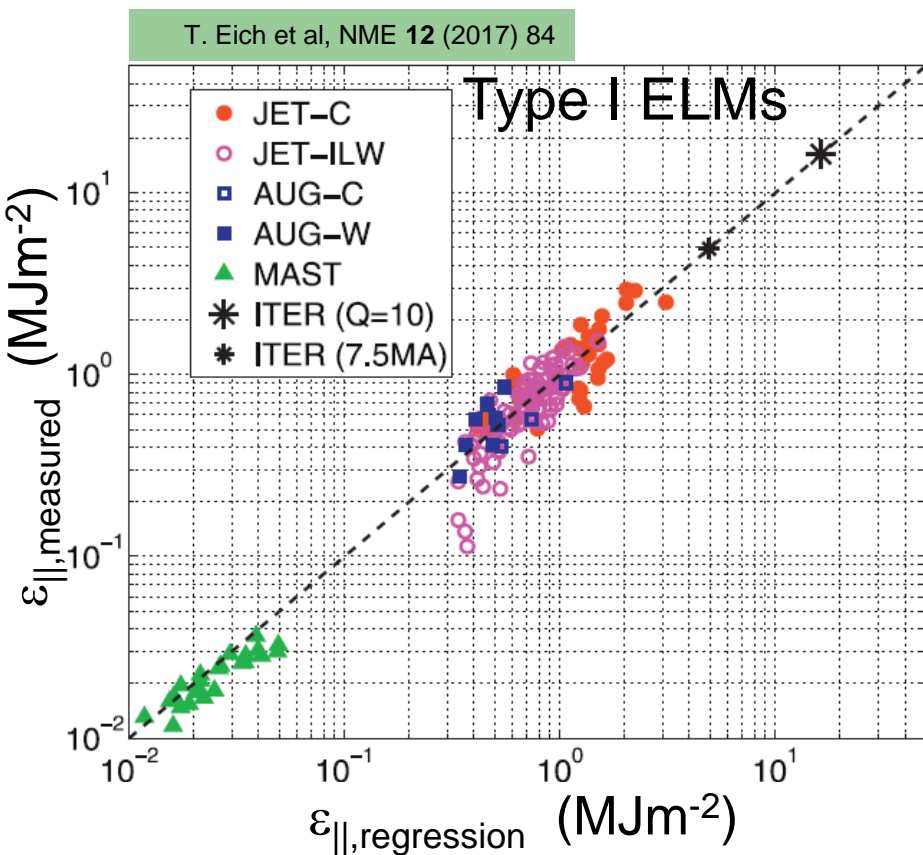
G. De Temmerman et al., PPCF 60 (2018) 044018, R. A. Pitts, NME submitted



- Add a few points from specific measurements (1 hour annealing) on ITER grade W materials¹
 - An idea of the range of uncertainty
 - Conclude that a reasonable max stationary heat flux could be $q_{\text{peak,target}} \lesssim 15 \text{ MWm}^{-2}$ for first ITER divertor to end of FPO

¹S. Panayotis et al., NME 12 (2017) 200

ELMs – what if suppression not possible?

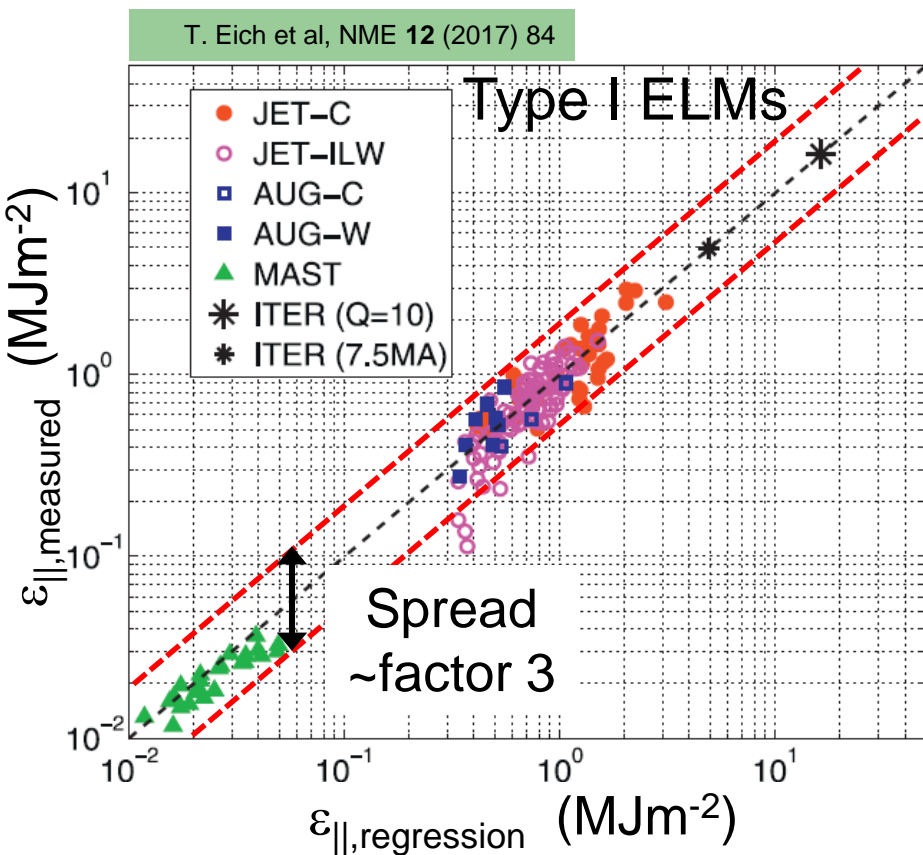


- Encouraging multi-device scaling for **outer target** peak parallel ELM energy density

$$\epsilon_{||,scaling} = 0.28 \frac{MJ}{m^2} n_{e,ped}^{0.75} T_{e,ped}^1 \Delta W_{ELM}^{0.52} R^1$$

- $\Delta W_{ELM} = W_{ELM} / W_{plasma}$
- $n_{e,ped}$, $T_{e,ped}$ values of n_e and T_e at the top of the H-mode pedestal

ELMs – what if suppression not possible?



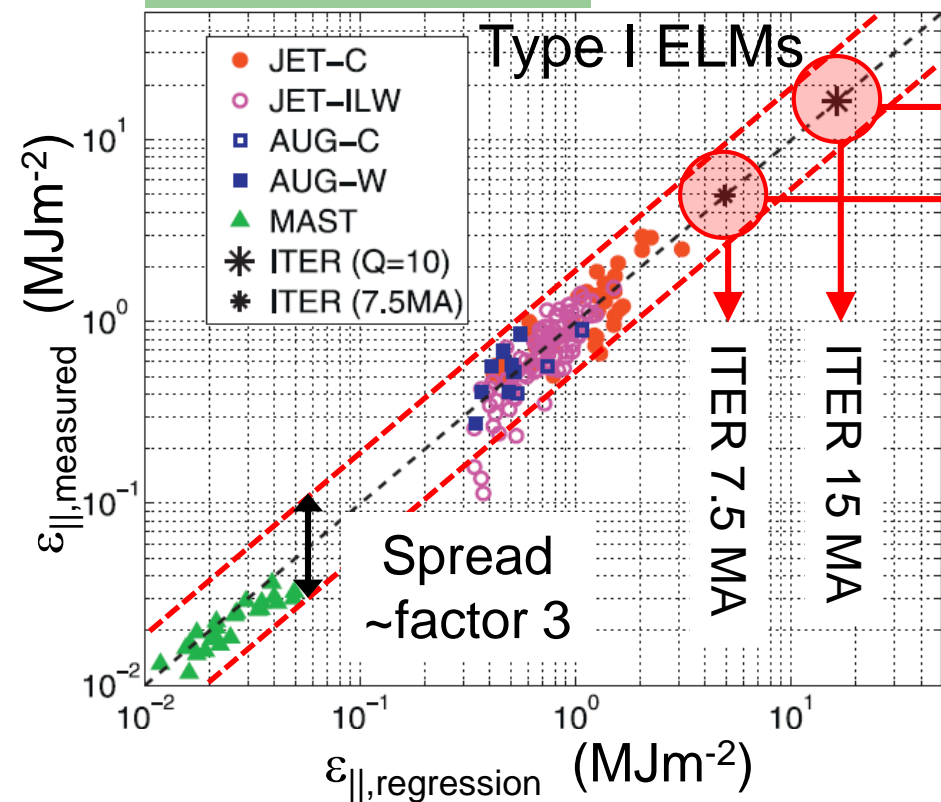
- Encouraging multi-device scaling for **outer target** peak parallel ELM energy density

$$\epsilon_{||,scaling} = 0.28 \frac{\text{MJ}}{\text{m}^2} n_{e,ped}^{0.75} T_{e,ped}^1 \Delta W_{ELM}^{0.52} R^1$$

- Parallel energy at targets dependent on pedestal top pressure and R
- Favourable for ITER at $Q_{DT} = 10$ compared to our previous scalings
- Lower bound of data matched by simple model: $\epsilon_{||} \approx 6\pi \cdot p_e R q_{edge}$ (pedestal plasma connects to the targets during the ELM)

ELMs – what if suppression not possible?

T. Eich et al, NME 12 (2017) 84

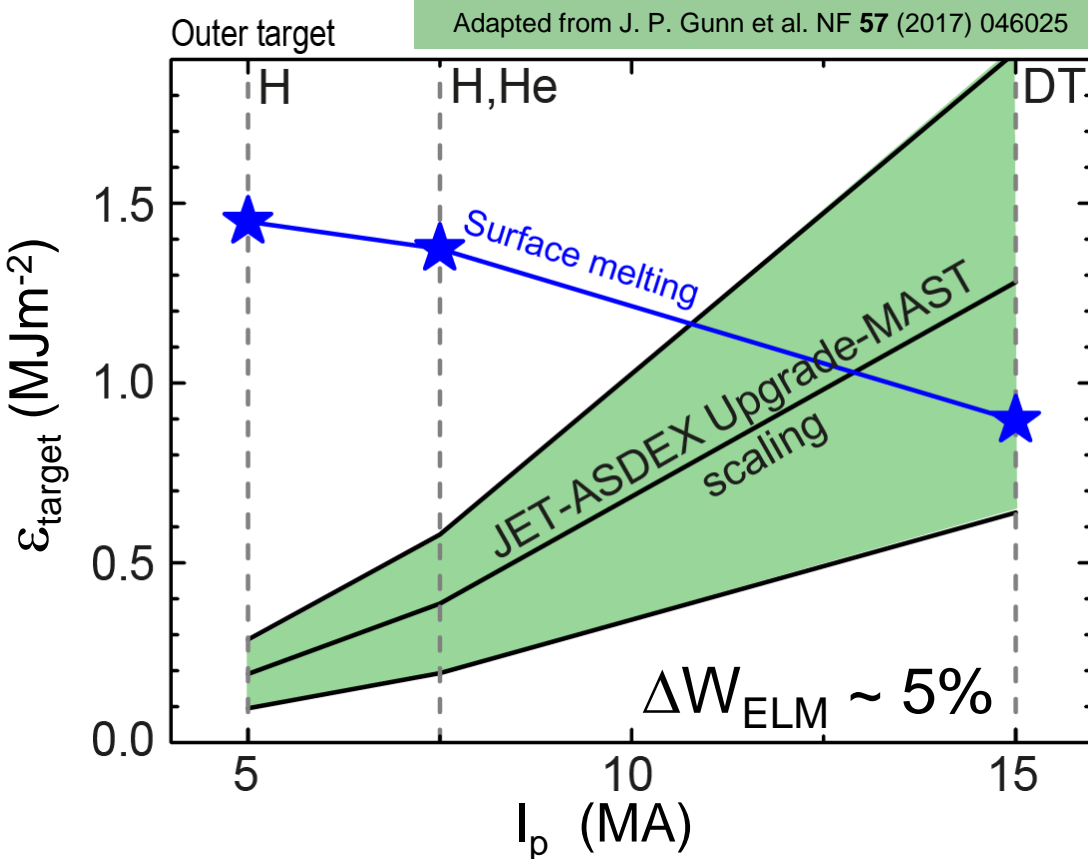


For $\Delta W_{ELM} \sim 5\%$

$\epsilon_{target} \sim 0.60 - 1.80 \text{ MJm}^{-2}$

$\epsilon_{target} \sim 0.18 - 0.54 \text{ MJm}^{-2}$

Natural Type I ELMs will still melt MB surface

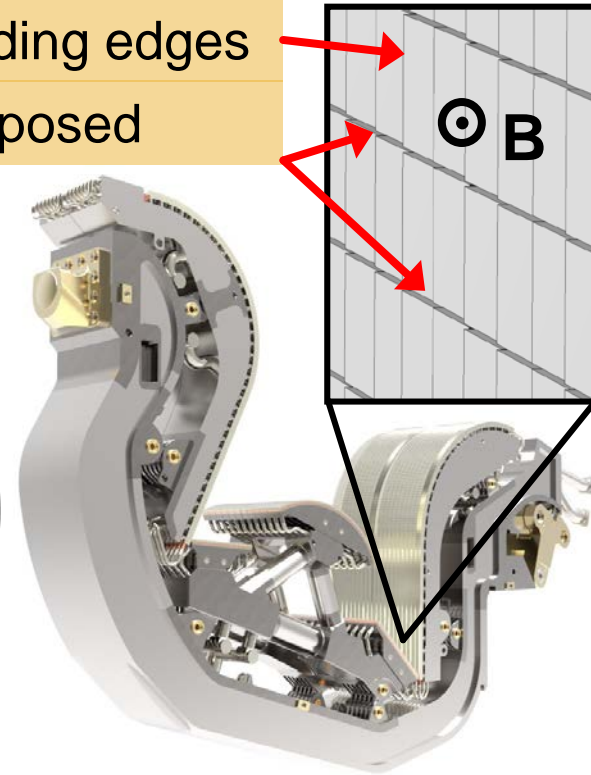
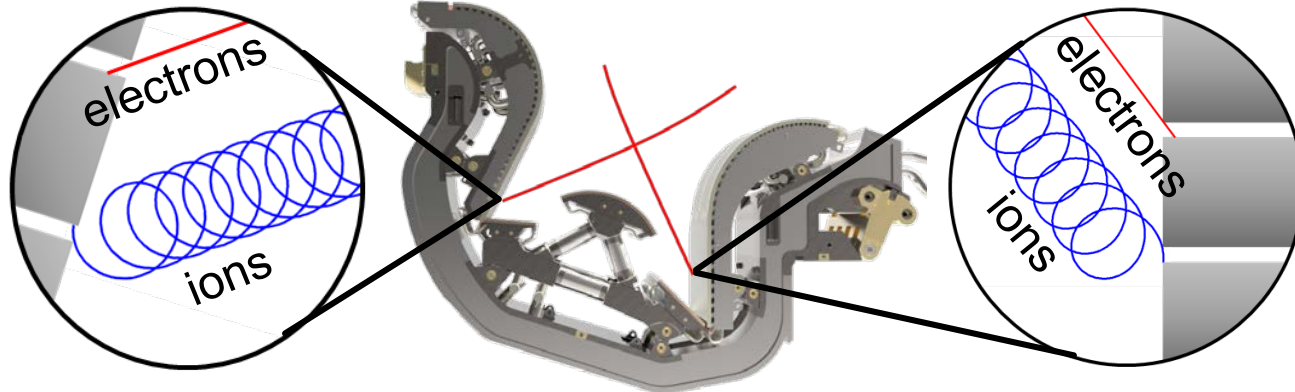


- Even with new scaling, monoblock surface melting will occur at $Q_{\text{DT}} = 10$
- Looks like reasonable margin to surface melting even for largest ΔW_{ELM}

Problem of tile gaps

Toroidal bevel protects poloidal leading edges
BUT long toroidal edges are still exposed

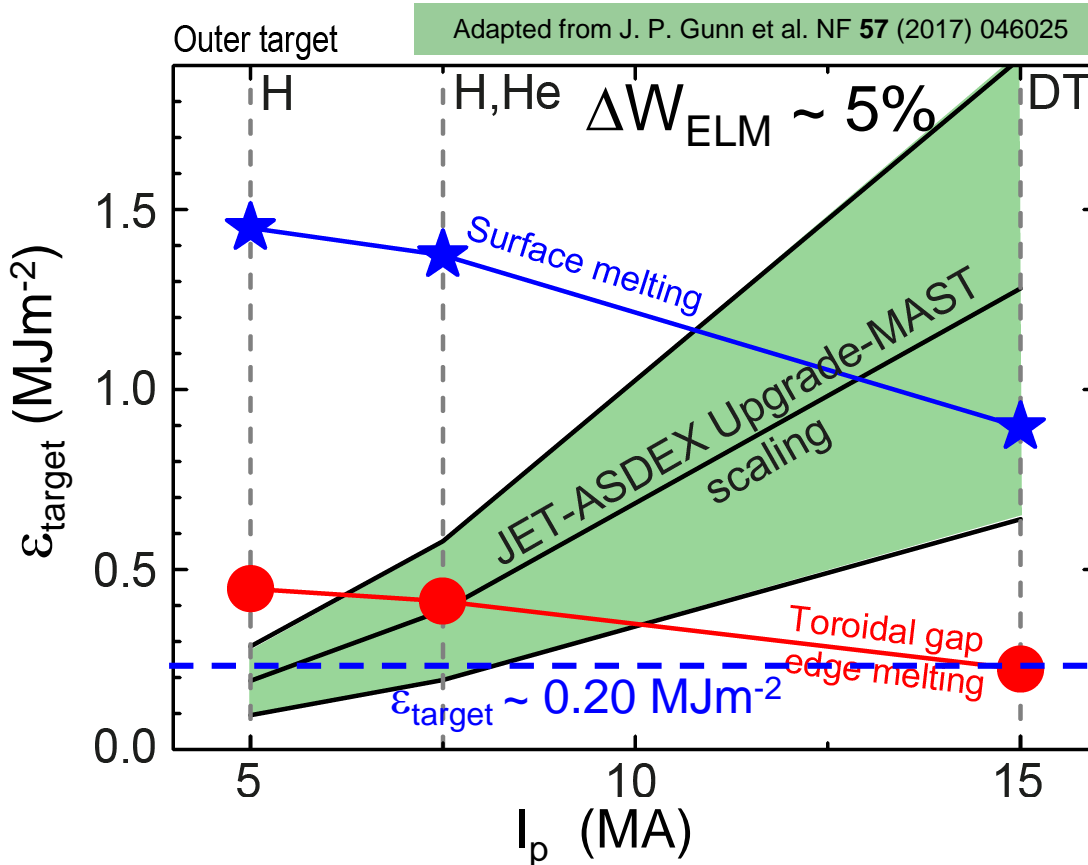
- ELM ions problematic due to large Larmor radii of particles arriving from pedestal region



- Toroidal gap (TG) loading really does occur¹
 - See talk by J. P. Gunn for more

¹R. Dejarnac, Nucl. Fus. **58** (2018) 066003

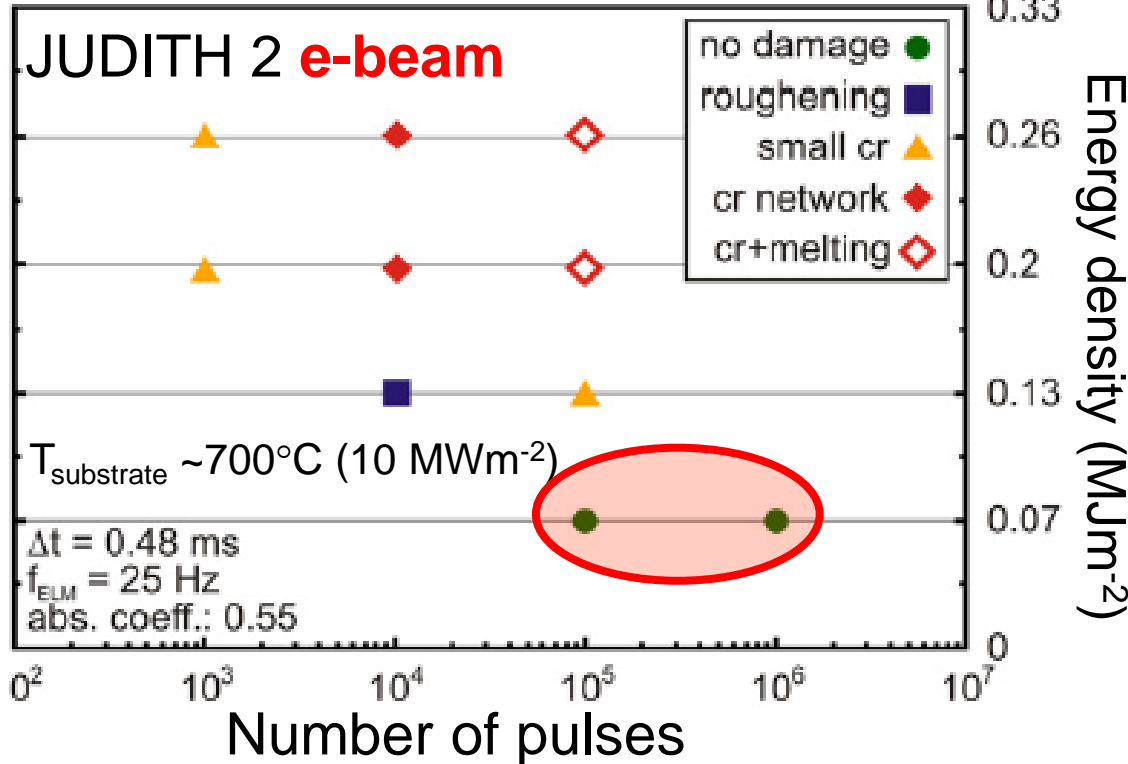
Now add toroidal gap melting



- To avoid toroidal gap edge melting:
 $\epsilon_{\text{target}} \lesssim 0.2 \text{ MJm}^{-2} \rightarrow$
 $\Delta W_{\text{ELM}}/W_{\text{plasma}} \lesssim 0.1\%$
 - Type I ELMs of this relative energy loss not found naturally

Surface cracking

T. Loewenhoff et al., Phys. Scr. T145 (2011) 014057



- Frequent thermal cycling can lead to W surface micro-cracking
 - Threshold for zero damage formation $\lesssim 0.1 \text{ MJm}^{-2}$ at high cycle number \rightarrow similar to toroidal gap edge melting
 - Micro-cracks may be initiators for larger macro-cracks

ELMs: what to do for ITER?

We don't know (yet) the consequences of repetitive ELM-induced monoblock toroidal gap melting

We don't know (yet) the consequences for fatigue-induced surface cracking under simultaneous plasma exposure

We know that ΔW_{ELM} must be kept below ~ 1 MJ to avoid monoblock top surface melting for $Q_{\text{DT}} = 10$ (15 MA, 5.3 T) $\rightarrow \sim 0.3\%$ of stored energy

Type I ELMs this small are not found naturally

So complete suppression is the only way to be sure.

But may also come with a price: see talks by Y. In, M. Fenstermacher, O. Schmitz

So that's it, ITER wall and divertor (enjoy your lunch!)

but come back to listen to H. Zohm to how all this looks for
the step beyond ITER. ..